

Magnetic field extrapolation in close binary stars

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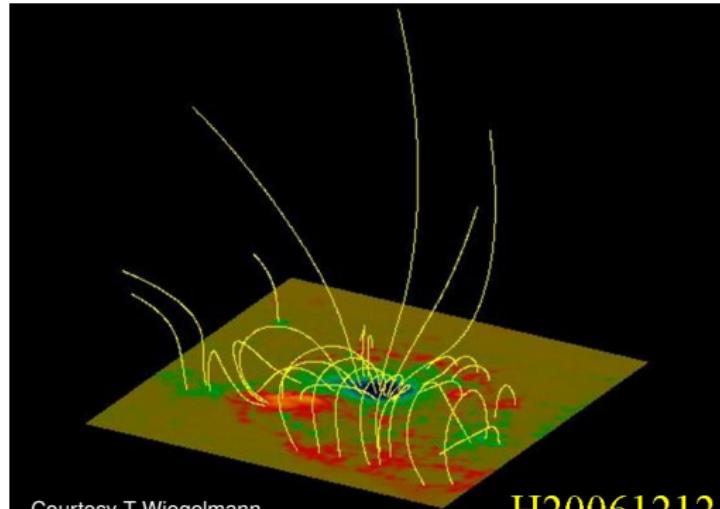
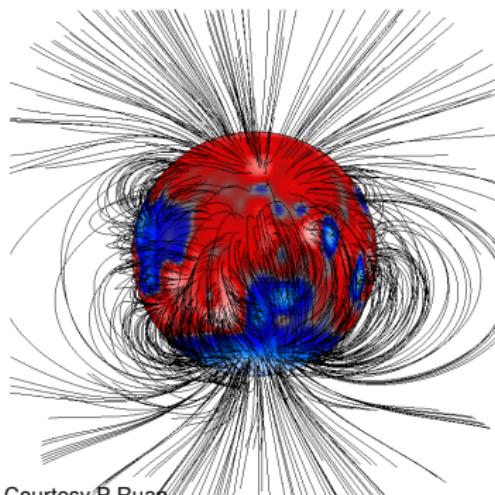
Binamics meeting, Vienna, Feb 2016



Solar magnetic fields

Coronal field extrapolation

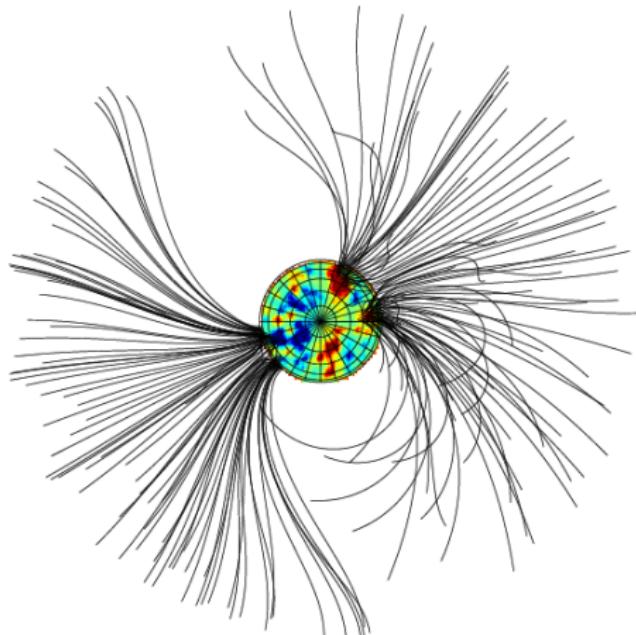
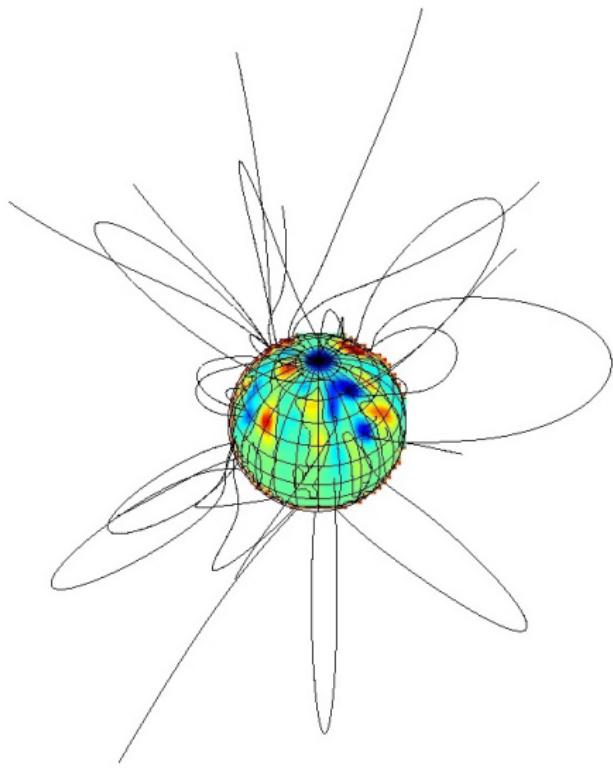
- observation of coronal magnetic field hardly possible
- ⇒ approximation techniques:
- **current-free** (potential) field approx.
 - **non-/linear force-free** field approx.
 - **magneto-hydrostatic** stratification



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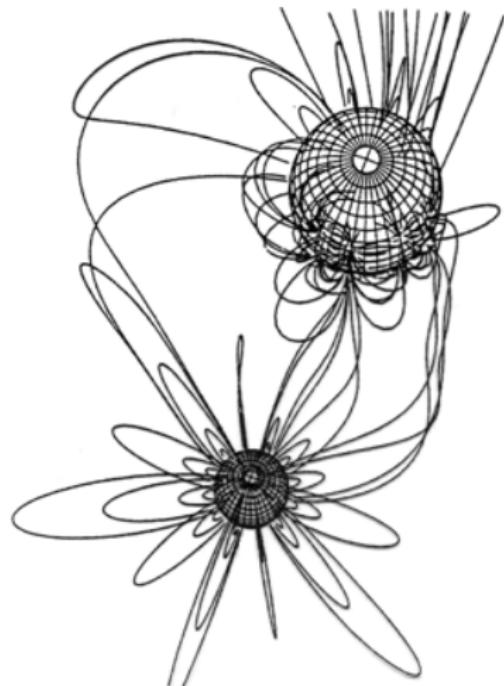
AB Dor

Potential field source surface (PFSS) approximation of coronal field structure (Jardine et al. 2002)



Joint magnetospheres of close star systems

Project overview



Uchida & Sakurai 1985

- targets:

- e.g.: RS CVn-/BY Dra-systems
- star-exoplanet systems
- circumbinary exoplanets
- hot binary stars
- ...

- objectives:

- structure of joint magnetosphere
- open/closed/inter-connecting magnetic flux
- magnetic interactions
- atmospheric structure
- observational signatures
- ...

Theoretical approach

PFSS approximation

- Assumption: **current-free** magnetic field

$$\nabla \times \mathbf{B} = 0 \quad \rightarrow \quad \mathbf{B} = -\nabla \Psi \quad \Psi: \text{scalar flux function}$$

- Solenoidal field:

$$\nabla \cdot \mathbf{B} = 0 \quad \rightarrow \quad \Delta \Psi = 0$$

→ **Laplace equation**

Theoretical approach

PFSS approximation

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→ **Laplace equation**

- General solution:

$$\Psi(\mathbf{r}) = \sum_{l,m} a_{lm} R_{lm}(\mathbf{r}) + b_{lm} I_{lm}(\mathbf{r})$$

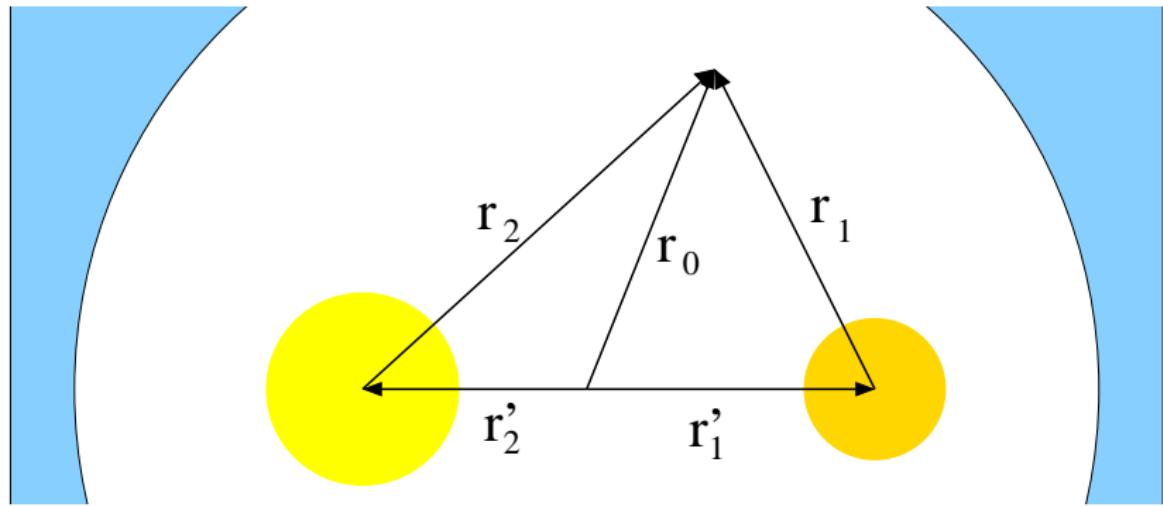
a_{lm}, b_{lm} expansion coefficients; Y_{lm} spherical harmonics

$R_{lm}(\mathbf{r}) = r^l Y_{lm}(\theta, \phi)$ **regular** solid spherical harmonics

$I_{lm}(\mathbf{r}) = \frac{1}{r^{l+1}} Y_{lm}(\theta, \phi)$ **irregular** solid spherical harmonics

Application to binary systems

Superposition of flux functions



$$\Psi(\mathbf{r}) = \Psi^{(0)}(\mathbf{r}_0) + \Psi^{(1)}(\mathbf{r}_1) + \Psi^{(2)}(\mathbf{r}_2)$$

- $\Psi^{(1)}(\mathbf{r}_1) = \mathbf{C}^{(1)} \cdot \mathbf{l}(\mathbf{r}_1)$ 'sources' inside **primary**
- $\Psi^{(2)}(\mathbf{r}_2) = \mathbf{C}^{(2)} \cdot \mathbf{l}(\mathbf{r}_2)$ 'sources' inside **secondary**
- $\Psi^{(0)}(\mathbf{r}_0) = \mathbf{C}^{(0)} \cdot \mathbf{R}(\mathbf{r}_0)$ 'mirror sources' outside **source surface**

Application to binary systems

Boundary conditions

- *Boundary conditions:*

- **Source surface:** Stellar winds drag magnetic field in radial direction

$$\Psi \Big|_{S_0} = \Psi^{(0)} \Big|_{S_0}(\mathbf{r}_0) + \Psi^{(1)} \Big|_{S_0}(\mathbf{r}_1) + \Psi^{(2)} \Big|_{S_0}(\mathbf{r}_2) = \text{const.}$$

- **Primary and secondary:** observed radial magnetic field maps

$$B_r^{(1)} = -\frac{\partial}{\partial r_1} (\Psi^{(0)}(\mathbf{r}_0) + \Psi^{(1)}(\mathbf{r}_1) + \Psi^{(2)}(\mathbf{r}_2)) \Big|_{S_1}$$

$$B_r^{(2)} = -\frac{\partial}{\partial r_2} (\Psi^{(0)}(\mathbf{r}_0) + \Psi^{(1)}(\mathbf{r}_1) + \Psi^{(2)}(\mathbf{r}_2)) \Big|_{S_2}$$

Application to binary systems

Boundary conditions

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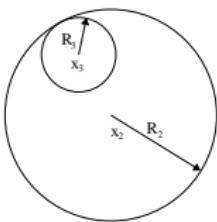
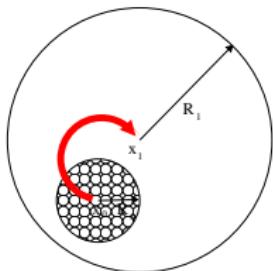
- **Primary and secondary:** observed radial magnetic field maps

$$B_r^{(1)} = -\frac{\partial}{\partial r_1} \left(\Psi^{(0)}(\mathbf{r}_0) + \Psi^{(1)}(\mathbf{r}_1) + \Psi^{(2)}(\mathbf{r}_2) \right) \Big|_{S_1}$$

$$B_r^{(2)} = -\frac{\partial}{\partial r_2} \left(\Psi^{(0)}(\mathbf{r}_0) + \Psi^{(1)}(\mathbf{r}_1) + \Psi^{(2)}(\mathbf{r}_2) \right) \Big|_{S_2}$$

- **Problem:** Express $\Psi^{(p)}$ in reference frame of sphere q
- **Solution:** **Multipole translation theorems**

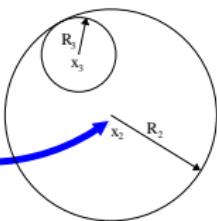
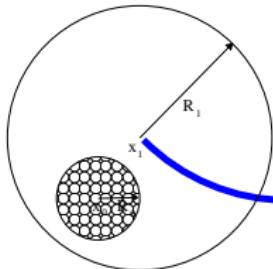
Multipole translation theorems



Irregular-irregular expansion: $|\mathbf{x} - \mathbf{x}_1| > |\mathbf{x}_1 - \mathbf{x}_0|$

$$I_{lm}(\mathbf{x} - \mathbf{x}_0) = \sum_{l',m'} (II)_{lm}^{l'm'} (\mathbf{x}_1 - \mathbf{x}_0) I_{l'm'}(\mathbf{x} - \mathbf{x}_1)$$

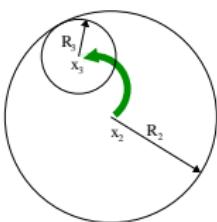
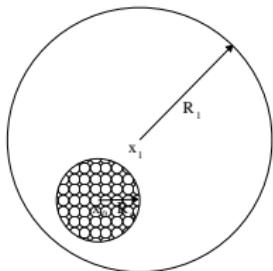
$$(II)_{lm}^{l'm'} (\mathbf{x}_1 - \mathbf{x}_0) = (-1)^{l'+m'-l-m} R_{l'-l}^{m-m'} (\mathbf{x}_1 - \mathbf{x}_0)$$



Irregular-regular expansion: $|\mathbf{x} - \mathbf{x}_2| < |\mathbf{x}_2 - \mathbf{x}_1|$

$$I_{lm}(\mathbf{x} - \mathbf{x}_1) = \sum_{l'm'} (II|R)_{lm}^{l'm'} (\mathbf{x}_2 - \mathbf{x}_1) R_{l'm'}(\mathbf{x} - \mathbf{x}_2)$$

$$(II|R)_{lm}^{l'm'} (\mathbf{x}_2 - \mathbf{x}_1) = (-1)^{l'-m'} I_{l+l'}^{m-m'} (\mathbf{x}_2 - \mathbf{x}_1)$$



Regular-regular expansion:

$$R_{lm}(\mathbf{x} - \mathbf{x}_2) = \sum_{l',m'} (R|R)_{lm}^{l'm'} (\mathbf{x}_3 - \mathbf{x}_2) R_{l'm'}(\mathbf{x} - \mathbf{x}_3)$$

$$(R|R)_{lm}^{l'm'} (\mathbf{x}_3 - \mathbf{x}_2) = R_{l-l'}^{m-m'} (\mathbf{x}_3 - \mathbf{x}_2)$$

Application to binary systems

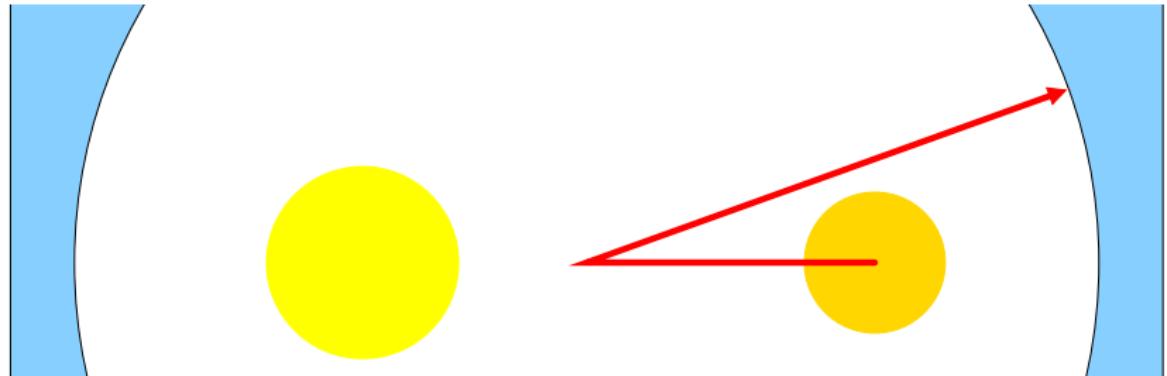
Example: Primary magnetic field



$$\begin{aligned}\mathbf{B}^{(1)} = & -\mathbf{C}^{(1)} \cdot \left[\mathcal{I}'_{a_1} - (I|I)^{(10)} \cdot \mathcal{I}_{a_0} \cdot \mathcal{R}_{a_0}^{-1} \cdot (R|R)^{(01)} \cdot \mathcal{R}'_{a_1} \right] \\ & - \mathbf{C}^{(2)} \cdot \left[(I|R)^{(21)} \cdot \mathcal{R}'_{a_1} - (I|I)^{(20)} \cdot \mathcal{I}_{a_0} \cdot \mathcal{R}_{a_0}^{-1} \cdot (R|R)^{(01)} \cdot \mathcal{R}'_{a_1} \right]\end{aligned}$$

Application to binary systems

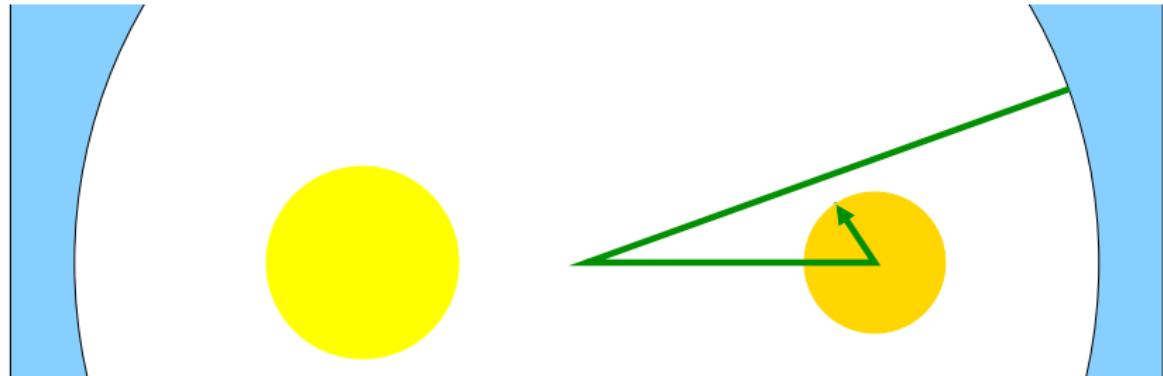
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Application to binary systems

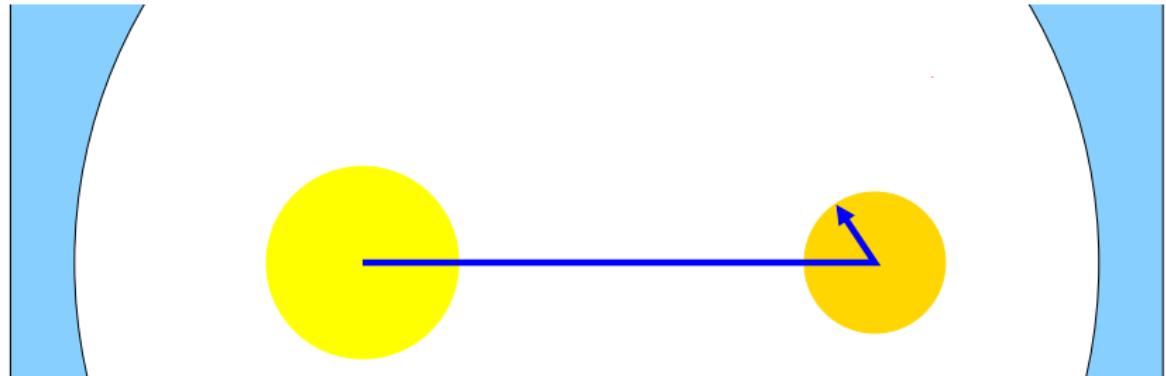
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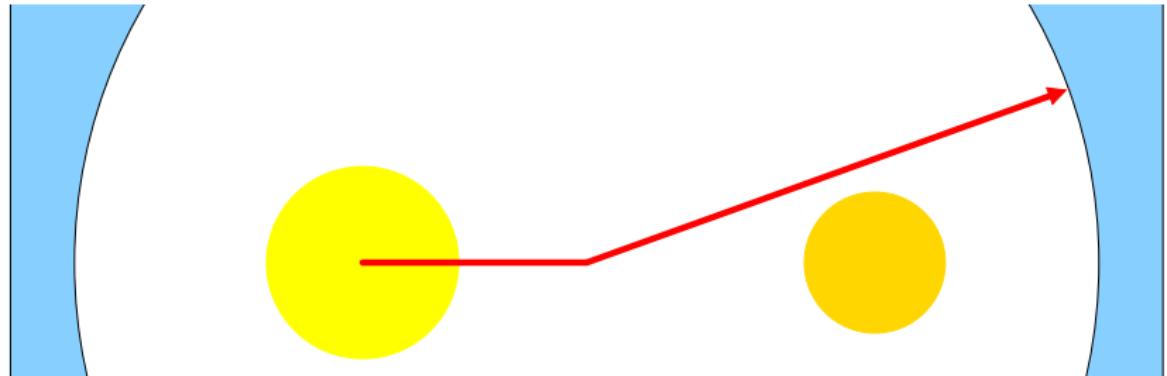
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Application to binary systems

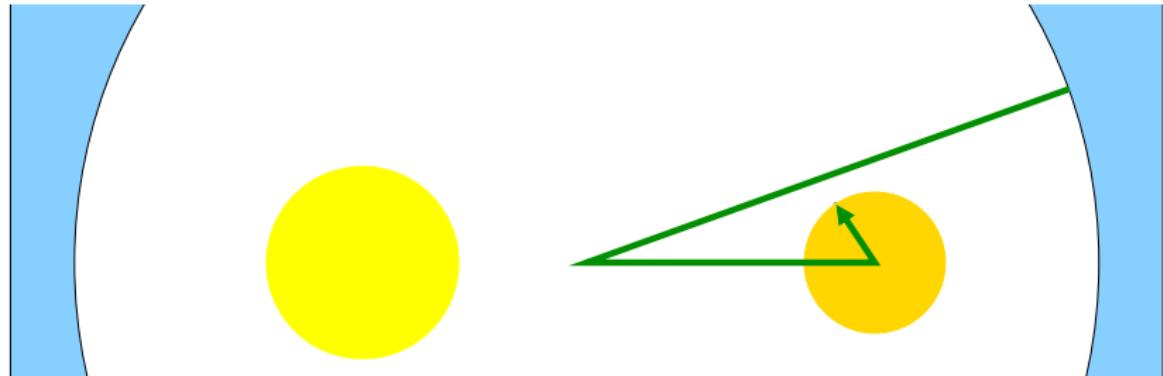
Example: Primary magnetic field



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Application to binary systems

Example: Primary magnetic field



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Application to binary systems

Linear system of equations

- binary system:

$$\begin{pmatrix} \mathbf{B}^{(1)} \\ \mathbf{B}^{(2)} \end{pmatrix} = \begin{pmatrix} \mathcal{D}^{(11)} & \mathcal{D}^{(12)} \\ \mathcal{D}^{(21)} & \mathcal{D}^{(22)} \end{pmatrix} \cdot \begin{pmatrix} \mathbf{C}^{(1)} \\ \mathbf{C}^{(2)} \end{pmatrix}$$

- magnetic maps $\mathbf{B}^{(1/2)}$ known \rightarrow solve for coefficients $\mathbf{C}^{(1/2)}$

Application to binary systems

Linear system of equations

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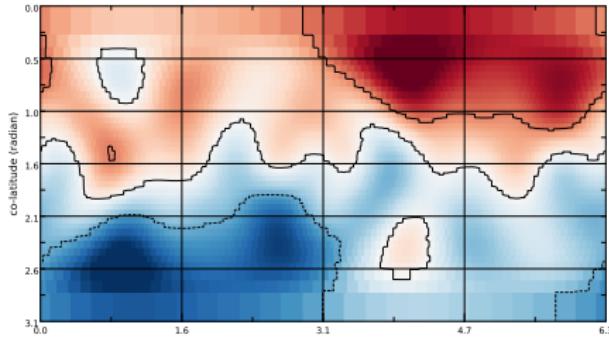
- magnetic maps $\mathbf{B}^{(1/2)}$ known \rightarrow solve for coefficients $\mathbf{C}^{(1/2)}$
- *post scriptum*: technique easily expandable to N objects

$$\begin{pmatrix} \mathbf{B}^{(1)} \\ \vdots \\ \mathbf{B}^{(j)} \\ \vdots \\ \mathbf{B}^{(N)} \end{pmatrix} = \begin{pmatrix} \mathcal{D}^{(11)} & \dots & \mathcal{D}^{(1j)} & \dots & \mathcal{D}^{(1N)} \\ \vdots & \ddots & \vdots & & \vdots \\ \mathcal{D}^{(j1)} & \dots & \mathcal{D}^{(jj)} & \dots & \mathcal{D}^{(jN)} \\ \vdots & & \vdots & \ddots & \vdots \\ \mathcal{D}^{(N1)} & \dots & \mathcal{D}^{(Nj)} & \dots & \mathcal{D}^{(NN)} \end{pmatrix} \cdot \begin{pmatrix} \mathbf{C}^{(1)} \\ \vdots \\ \mathbf{C}^{(j)} \\ \vdots \\ \mathbf{C}^{(N)} \end{pmatrix}$$

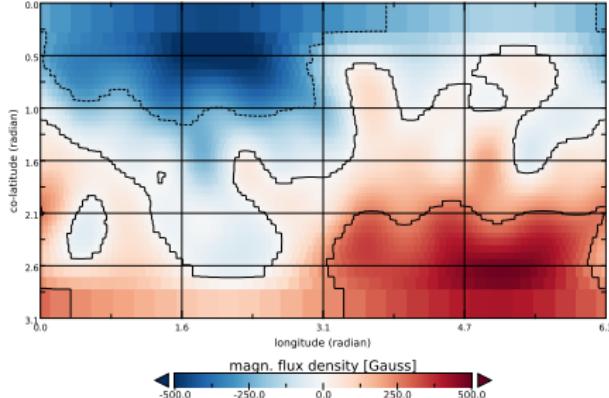
Example: YY Gem

Magnetic surface maps (Donati et al. 2011)

YY Gem, primary, radial magnetic field



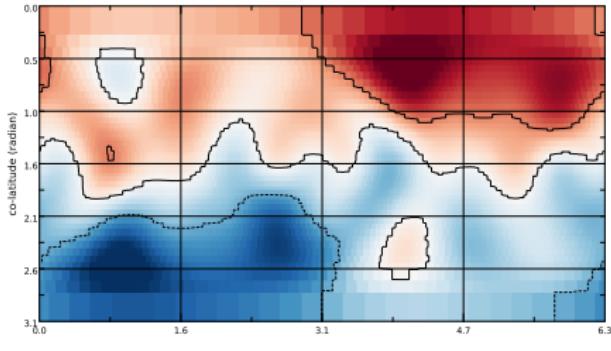
YY Gem, secondary, radial magnetic field



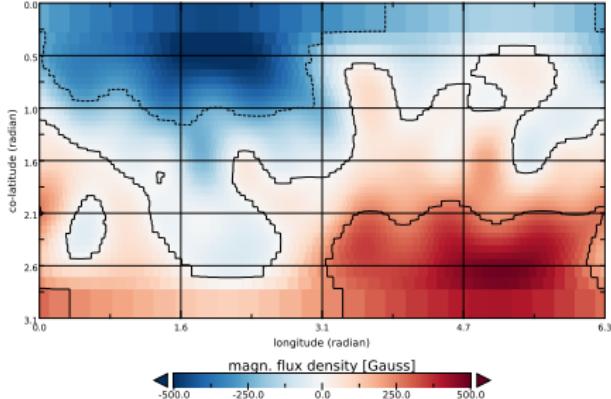
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Magnetic surface maps (Donati et al. 2011)

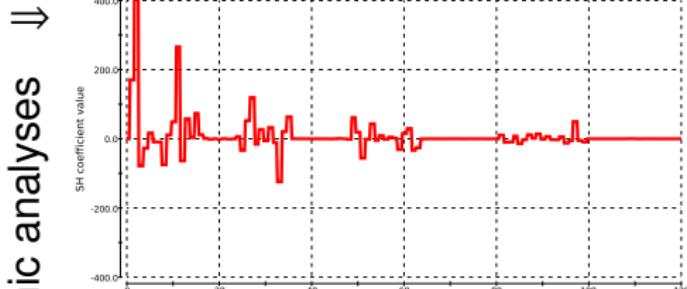
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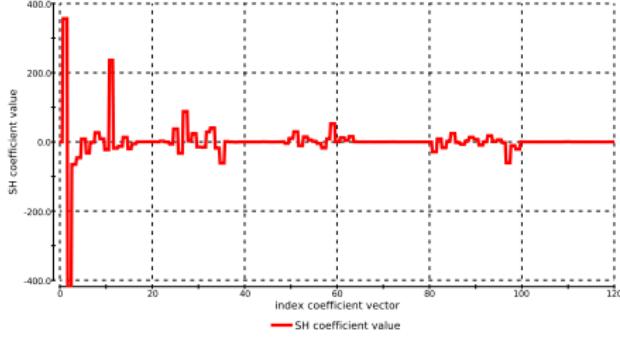


YY Gem, primary, magn. SH coefficient vector



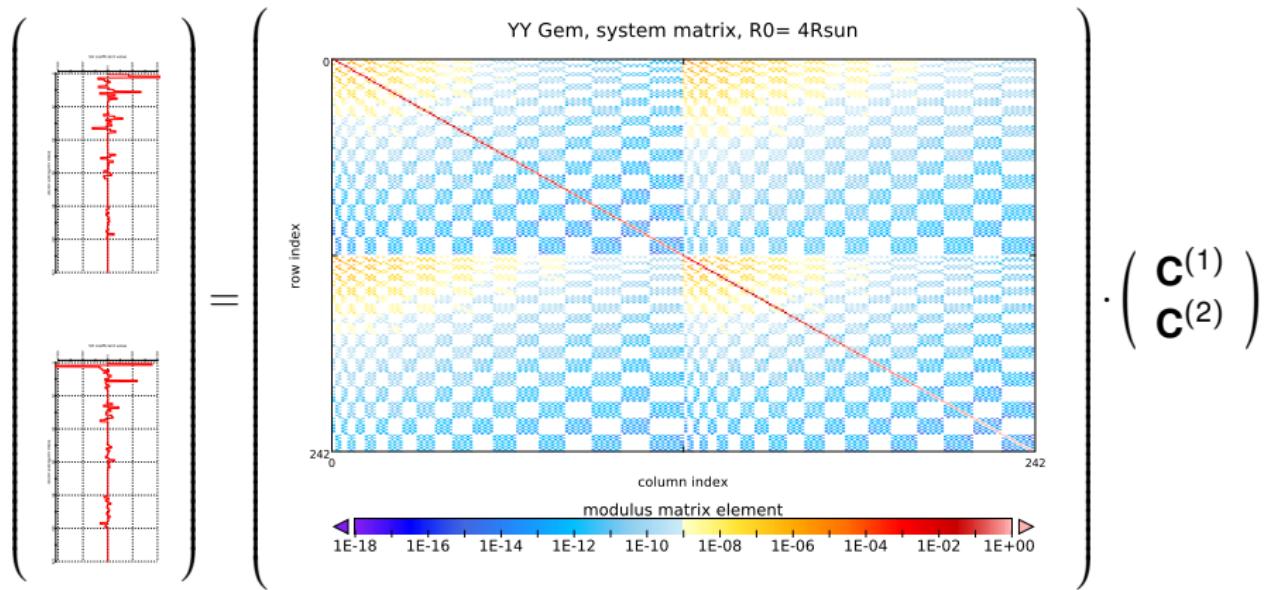
↓ spherical harmonic analyses

YY Gem, secondary, magn. SH coefficient vector



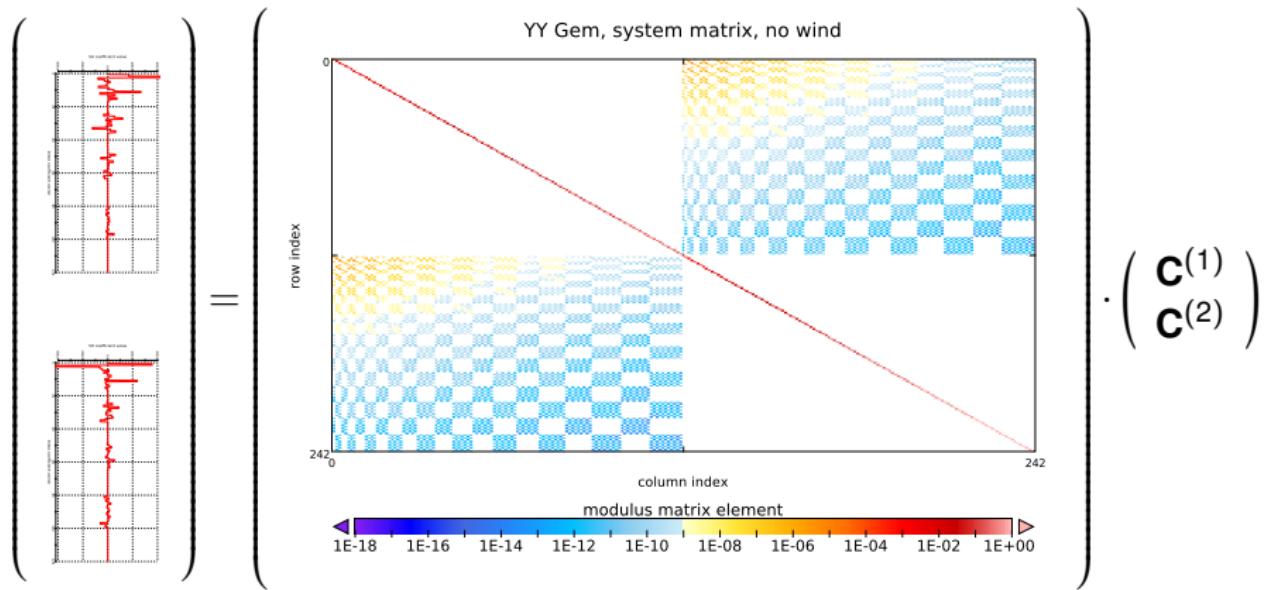
Example: YY Gem

Linear system of equations ($R_0 = 4R_\odot$)



Example: YY Gem

Linear system of equations ($R_0 \rightarrow \infty$)



Example: YY Gem

Joint magnetosphere – source surface radius $R_0 \gg d_{\text{YYGem}}$

