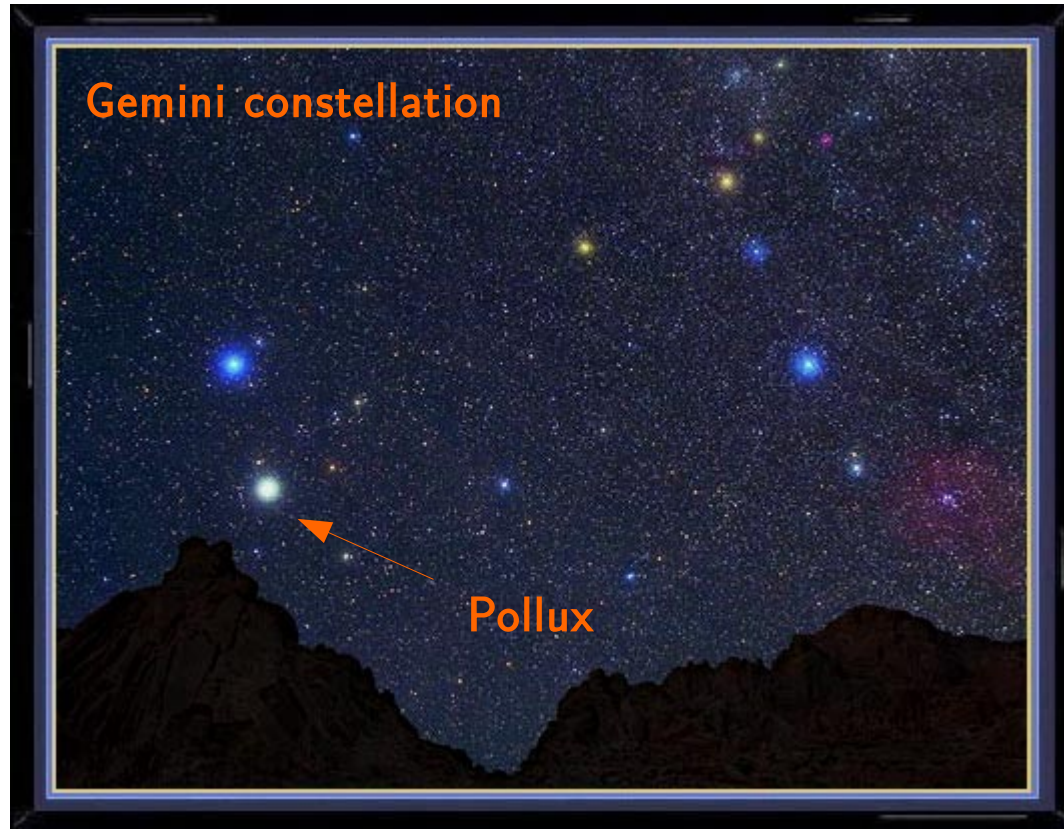
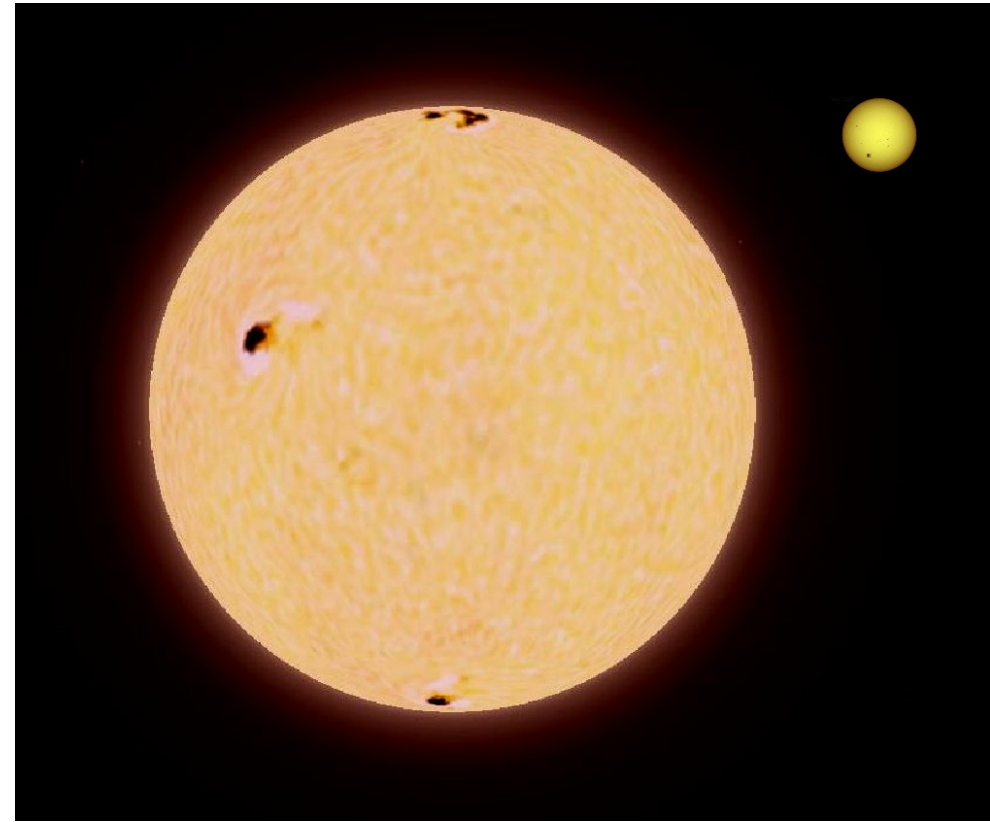


Magnetism of red giants stars

The case of the slow rotator, weakly magnetic giant POLLUX



© Bill and Sally Fletcher



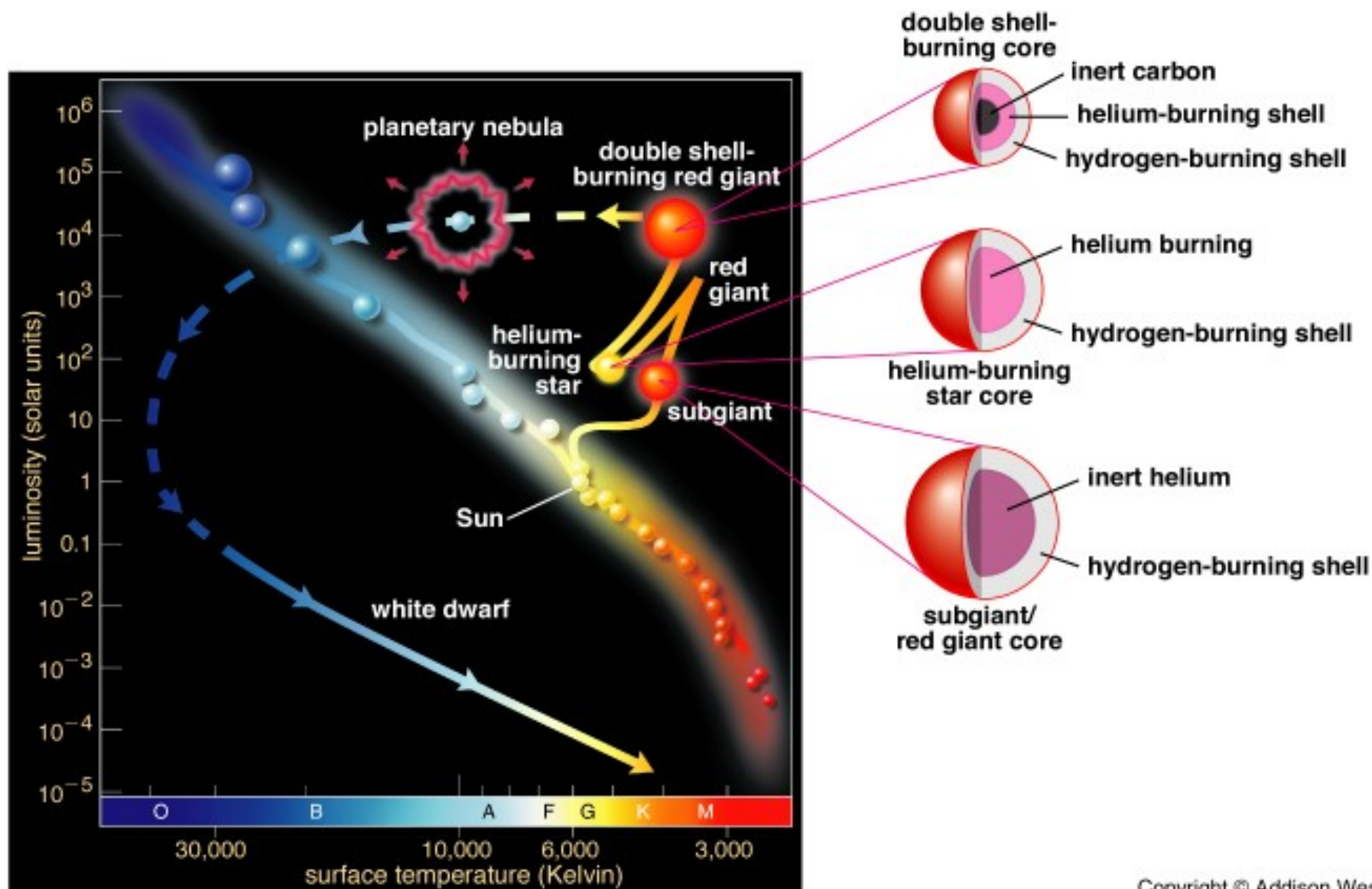
A. Palacios & A.-S. Brun
University of Montpellier / CEA-Saclay

Outline

- Some basics about red giants
- Magnetism of single giants (Aurière/Konstantinova-Antova group)
- The case of the weak-field giant Pollux
- 3D MHD simulations of the CE of Pollux
- Evolution of hydrodynamical quantities
- Evolution of magnetic field
- Impact of parameters : preliminary results

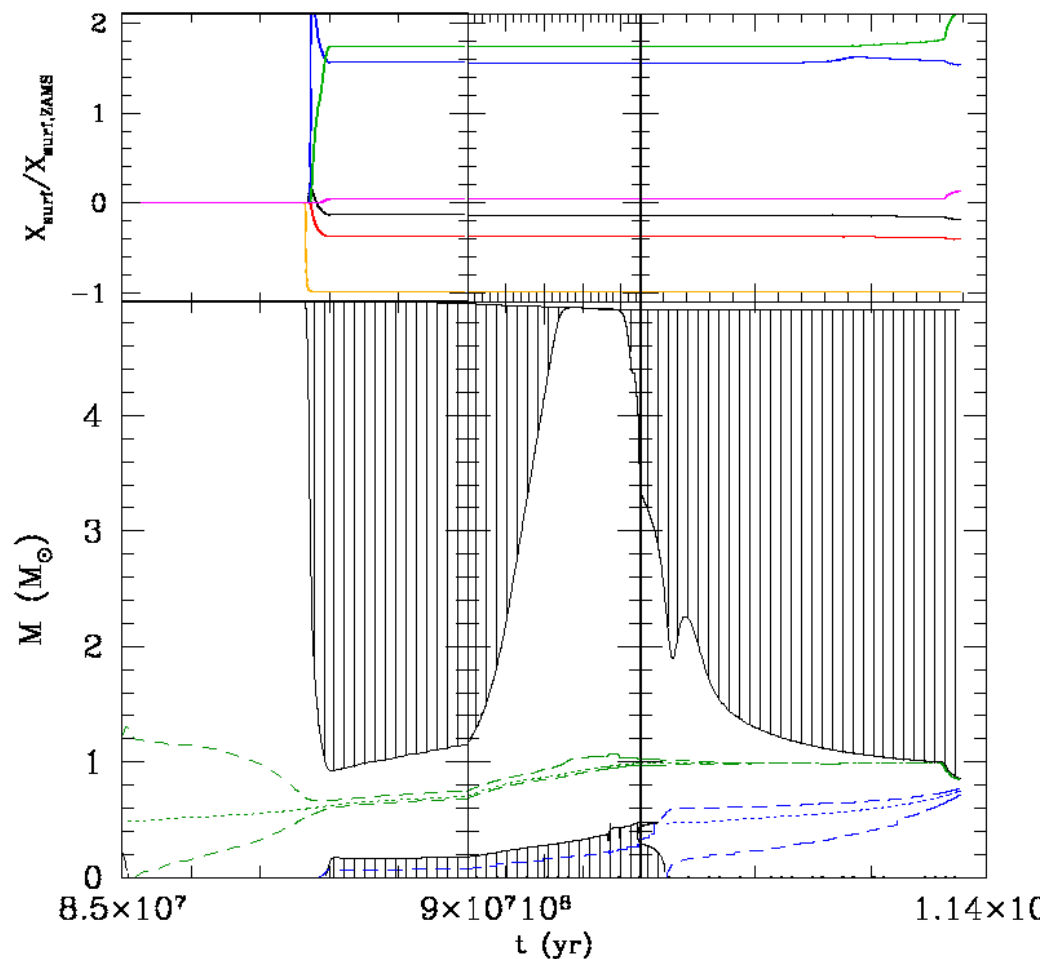
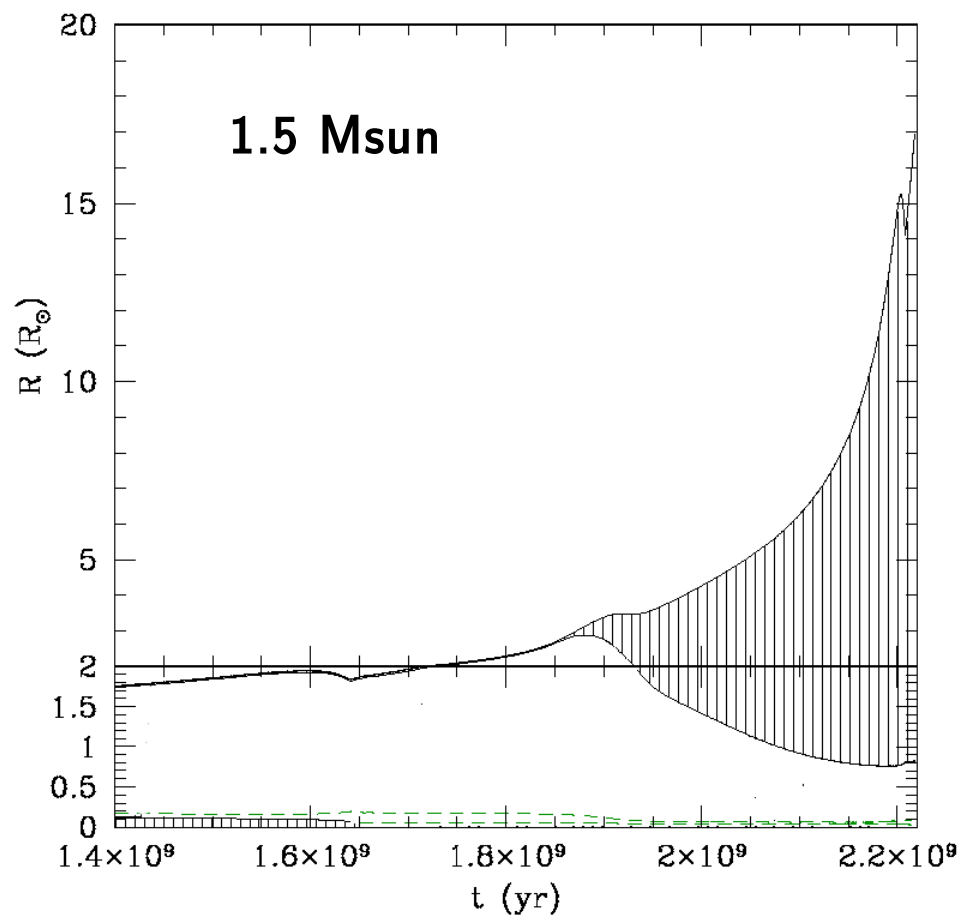
Red giant branch stars

Changing structure along evolutionary path : good laboratories to study dynamo



Internal structure of red giant stars

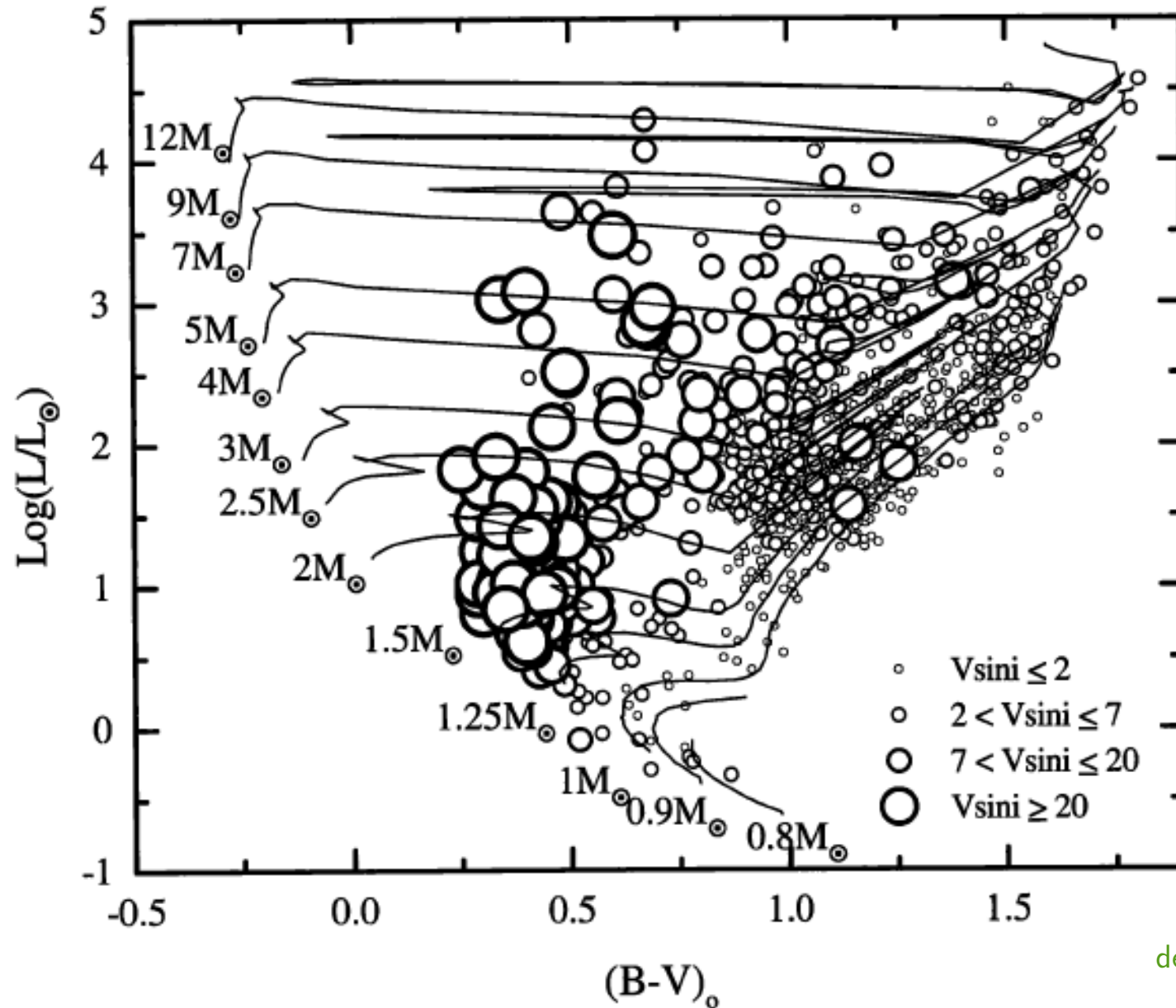
Very extended convective envelopes, low density, essentially slow rotation (challenged by asteroseismology)



5 Msun

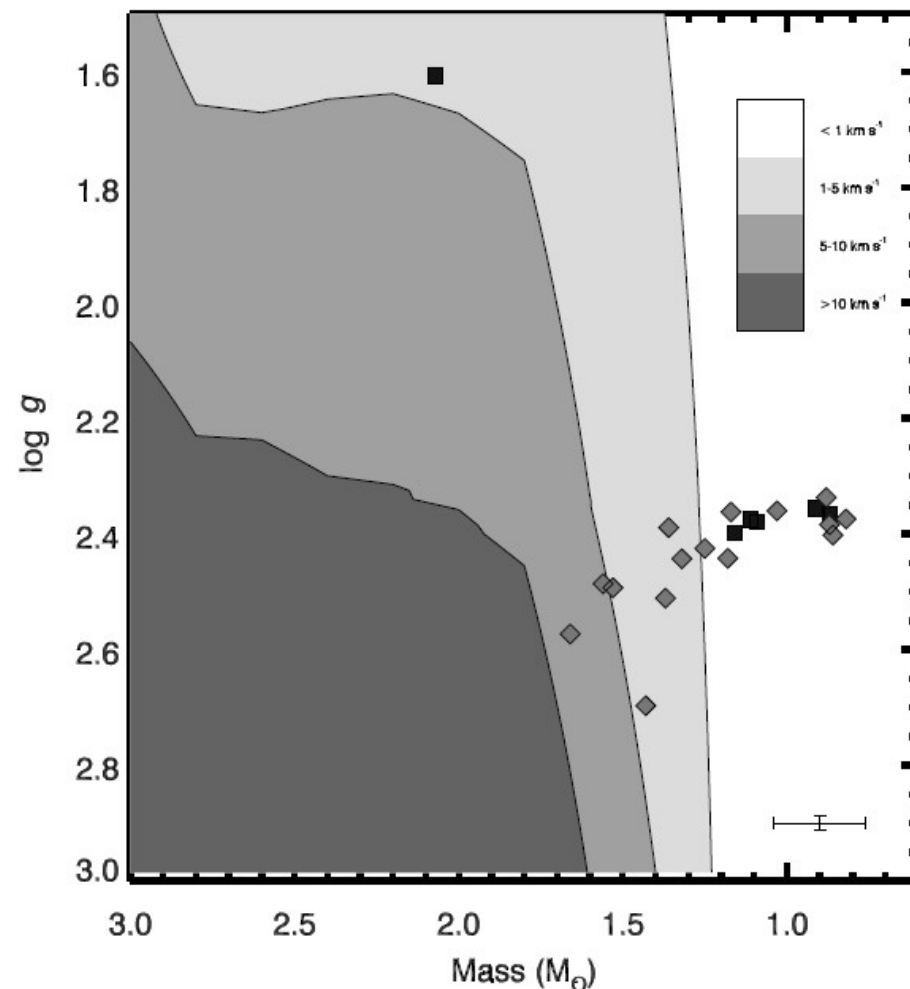
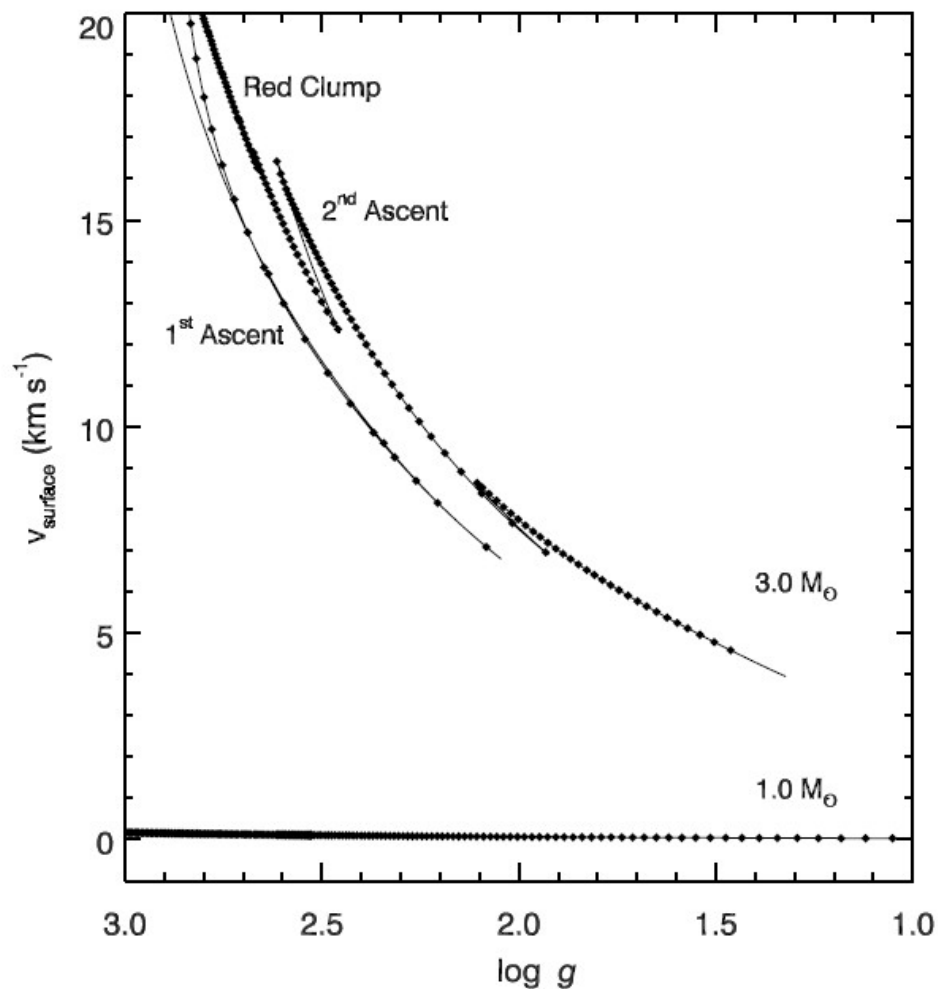
Observed rotation for red giant stars from spectroscopy

Essentially slow rotation on RGB/AGB with some exceptions



Predicted rotation of red giant stars

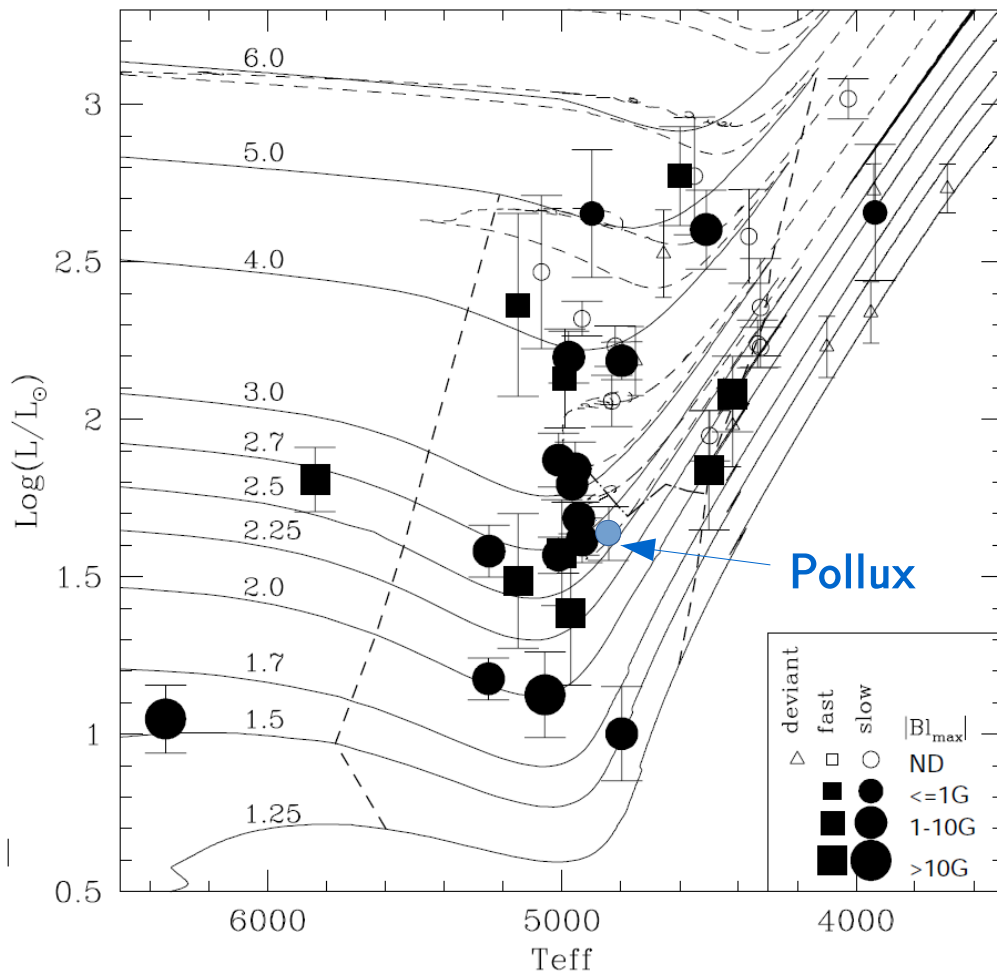
Surface velocity expected to decrease as the stars ascend the RGB (then AGB) with possible acceleration at clump due to overall contraction during CheB phase.



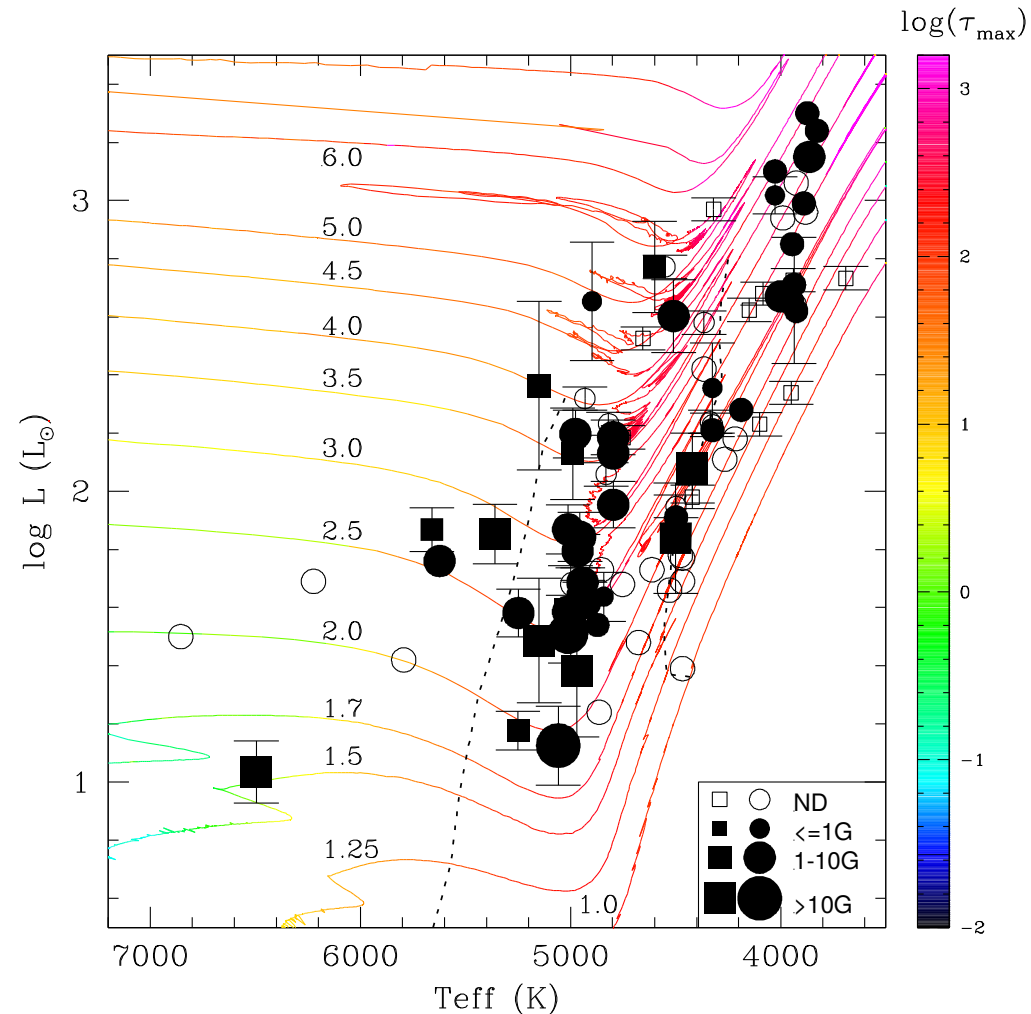
Magnetism of RGB stars

Two large spectropolarimetric surveys of single red giants in Stokes I and V by french-bulgarian team.

Selection of active and fast rotating red giants



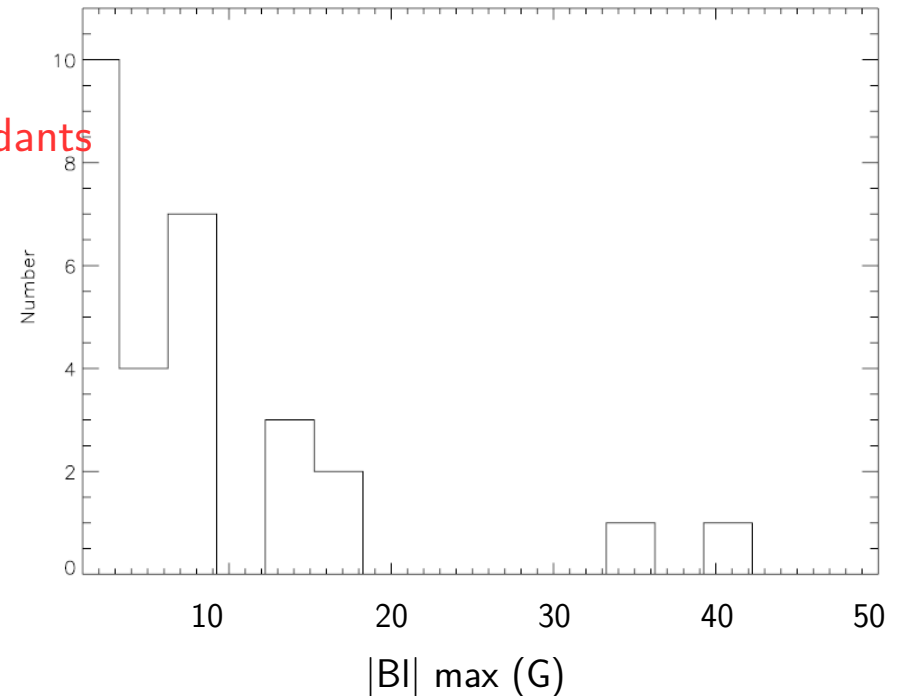
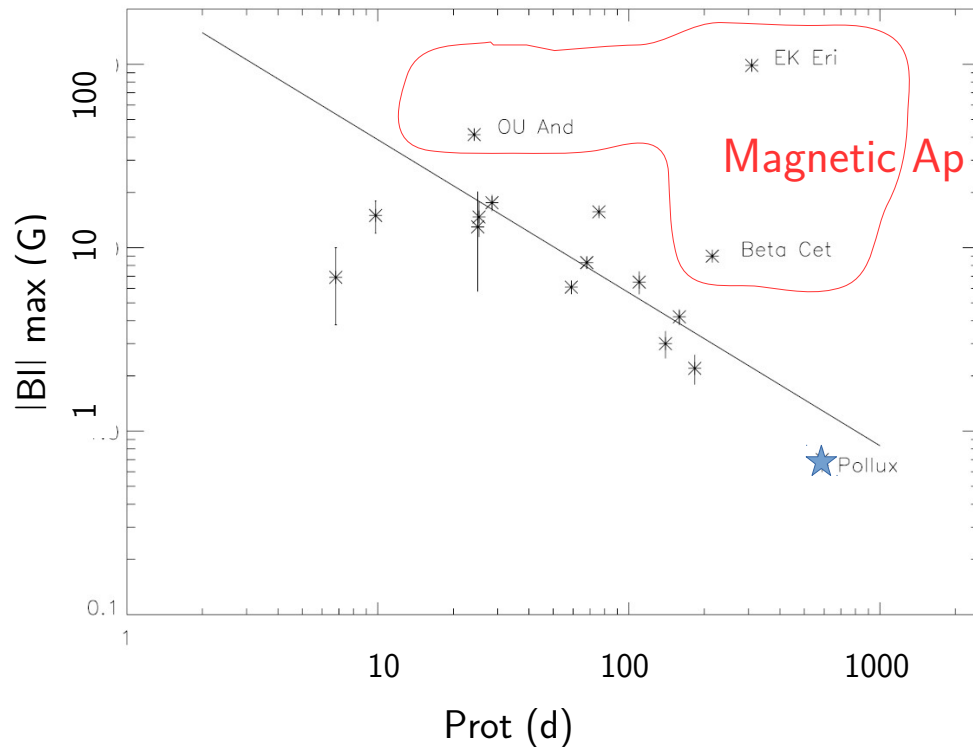
all single G, K and M giants up to 4 mag in the Solar vicinity observable from Pic du Midi



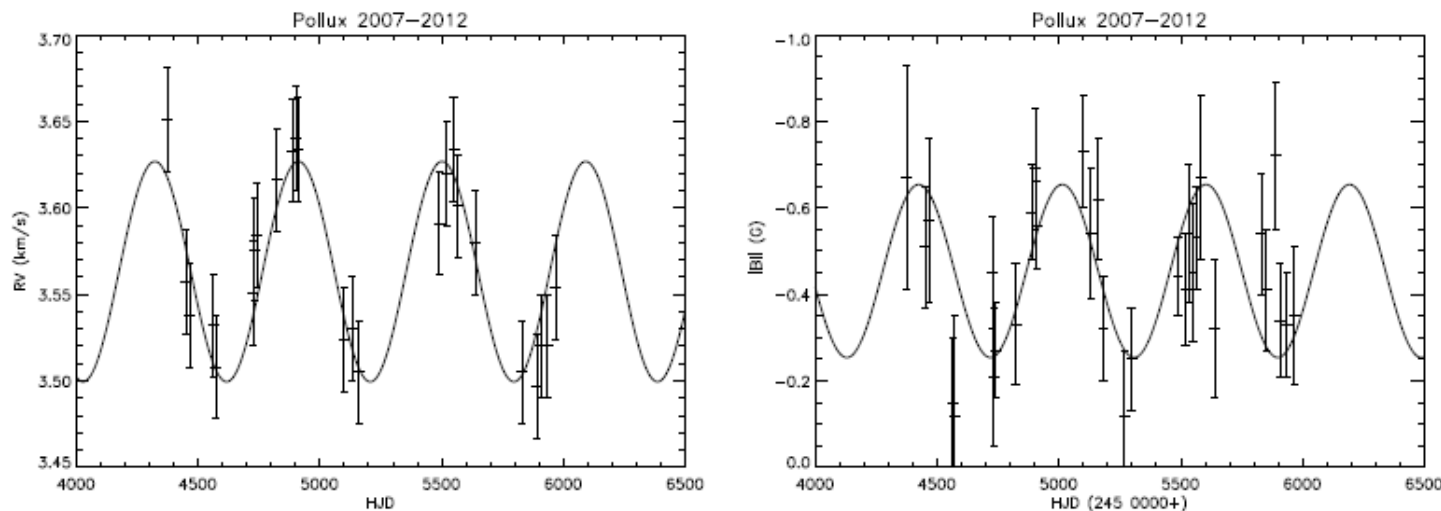
Magnetism of RGB stars – Main results from Aurière+ 2015

- ★ 29 of the 48 giants in the unbiased sample Zeeman detected in Stokes V
- ★ 23 out of 24 active giants Zeeman detected

- ★ Most of detected giants in magnetic strips where convective turnover timescale is maximum
- ★ The longer the rotation period the smaller the large scale longitudinal field



The case of the weak-field magnetic giant Pollux



Aurière et al. 2014, IAUS Conf. Proc. 302

Figure 1. Variations of RV (left plot) and of B_l (right plot) with HJD (245 0000+) in 2007-2012. A sinusoid with $P=589.64$ d is fitted for each parameter.

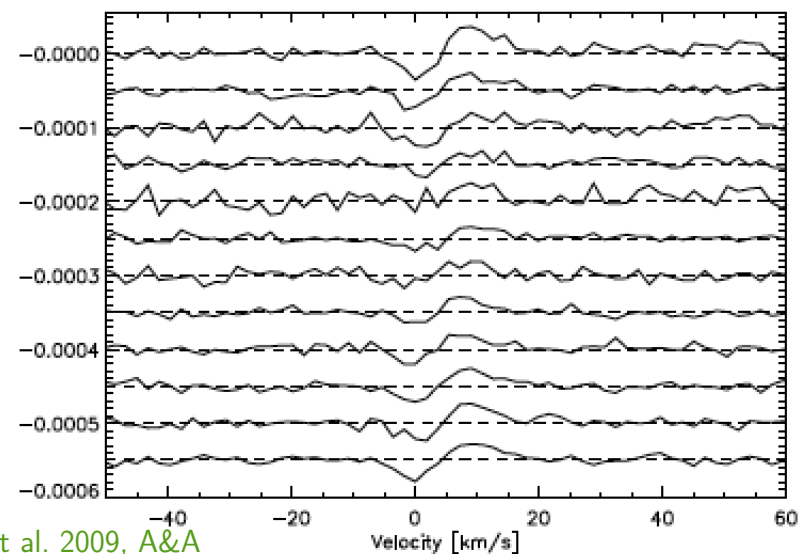
Table 2. Magnetic strength of red giants with available magnetic field models from ZDI.

Aurière et al. 2015, A&A

HD	Name	$v \sin i$ km s^{-1}	P_{rot} day	$ B_l _{\text{max}}$ G	B_{mean} G	Ref.
11812	31 Com	67	6.8	9.9	32	(1)
33798	V390 Aur	29	9.8	13	26	(2)
223460	OU And	21.5	24.2	36	68	(1)
112989	37 Com	11	111	6.5	10.8	(3)
9746	OP And	8.7	76	16	15	(4)
4128	β Cet	5.8	215	8	10	(5)
27536	EK Eri	1	308.8	99	94	(6)
62509	Pollux	2.8	590	0.7	0.6	(7)

Notes. References for B_{mean} : (1) Borisova et al. (in prep.); (2) Konstantinova-Antova et al. (2012); (3) Tsvetkova et al. (in prep.); (4) Konstantinova-Antova (in prep.); (5) Tsvetkova et al. (2013); (6) Aurière et al. (2011); (7) Aurière et al. (2014 and in prep.).

Very weak longitudinal field revealed by spectropolarimetry and long rotation period. ZDI reconstruction favours a strong dipole.



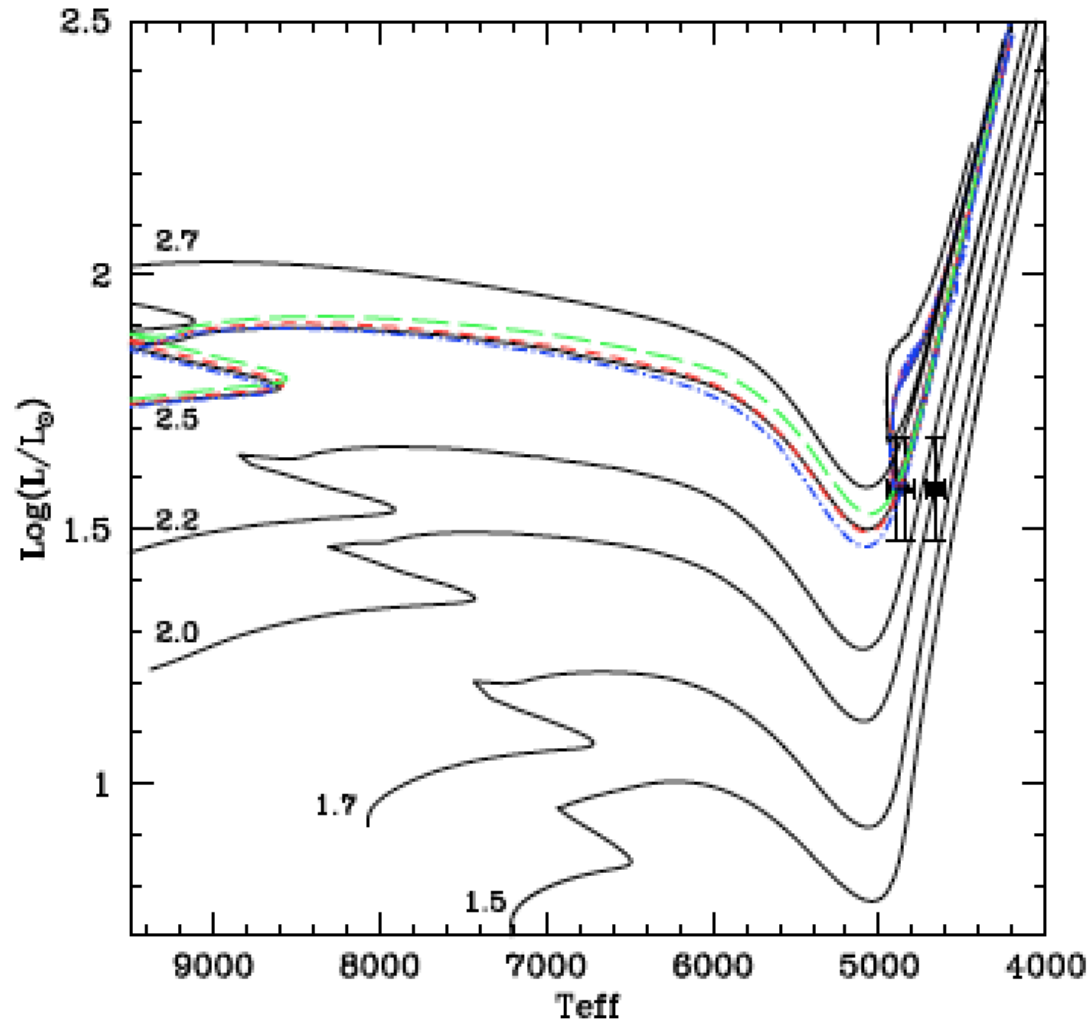
Aurière et al. 2009, A&A

Pollux – β Geminorum

RGB star K0III @ 10.3 pc of the Sun

$M_* = 2.5 M_{\text{sol}}$ $L_* = 40 L_{\text{sol}}$ $R_* = 9 R_{\text{sol}}$

Progenitor: $2.5 M_{\text{sol}}$ A2V star



ASH simulations

3D MHD simulations in the anelastic approximation using the ASH (Anelastic Spherical Harmonics) code (*Clune et al. 1999, Brun et al. 2004, Featherstone et al. 2013*)

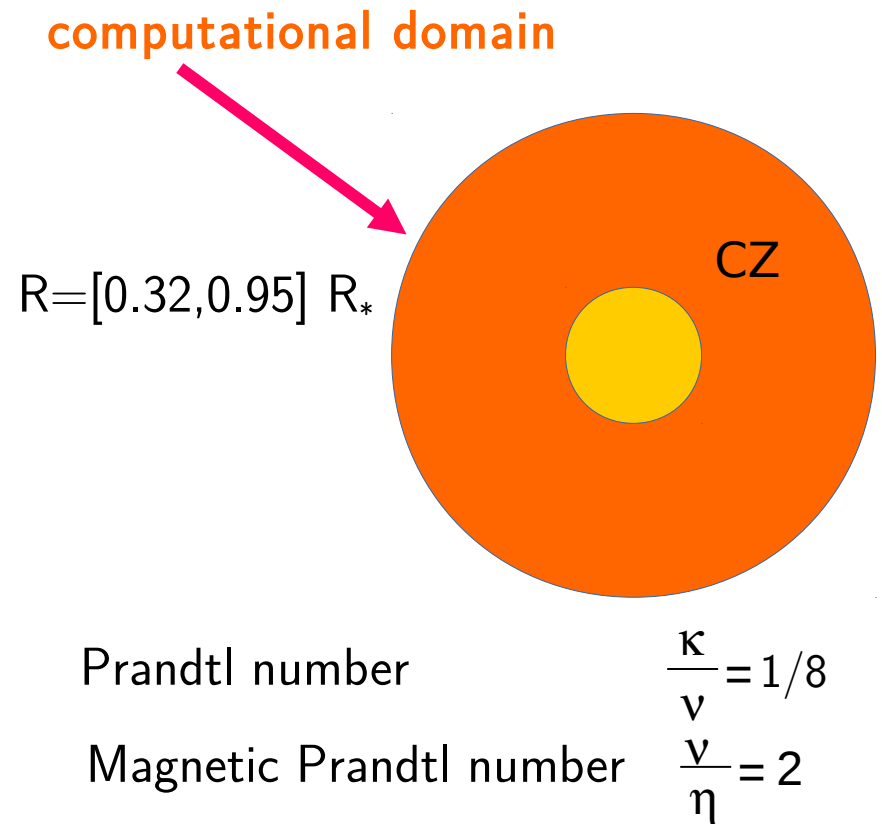
MHD runs based on converged hydrodynamical runs in which a magnetic multipole seed is included

- ⊗ $L \approx$ constant in model
- ⊗ $\rho \sim [0.02; 10^{-5}] \text{ g/cm}^3$ or 7.4 Np
- ⊗ $\Omega = \Omega_{\odot}/20 = 1.23 \cdot 10^{-7} \text{ rad s}^{-1}$
or $P_{\text{rot}} = 590 \text{ days}$
- ⊗ Resolution : $N_r \times N_{\theta} \times N_{\varphi} = 513 \times 512 \times 1024$

Rigid stress free boundaries, Fixed Flux
Perfect conductor (bottom) &
potential field (top)

Case A : initial high degree multipole $l = 10, m = 7$, $Re = 238$, $Ro = 2.35$

Case B : initial dipole $l = 1, m = 0$



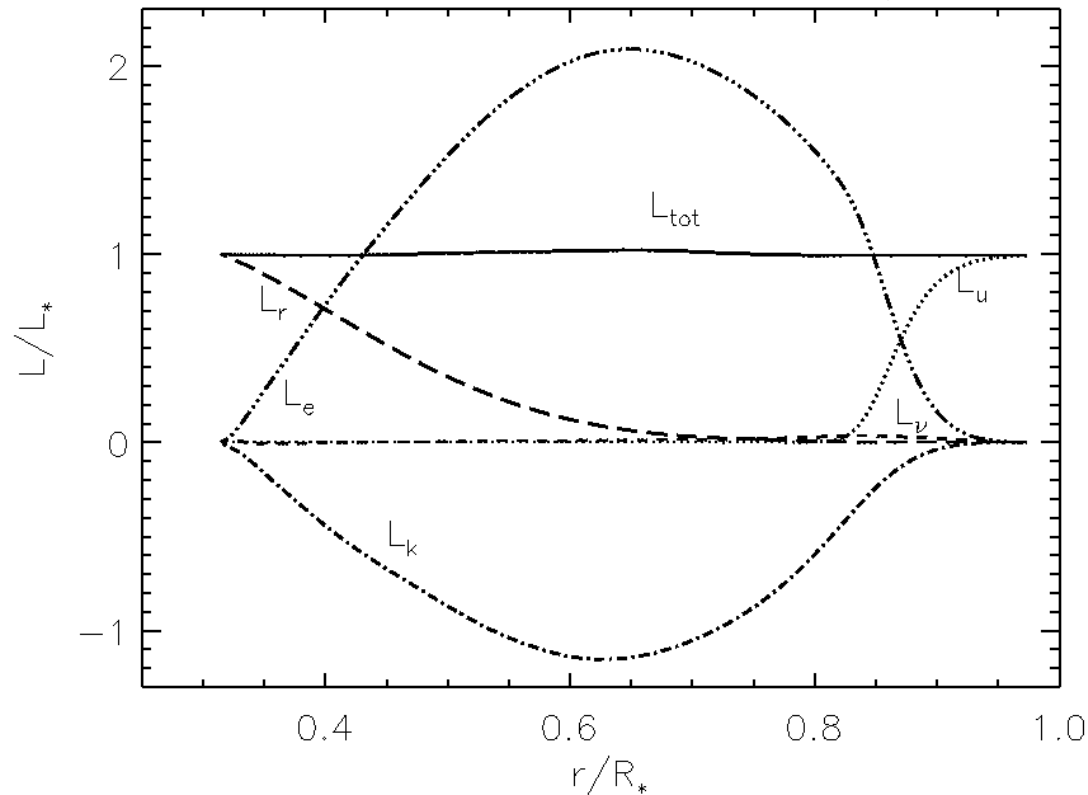
Energy budget

Convective luminosity $> L_*$ to compensate the negative kinetic luminosity

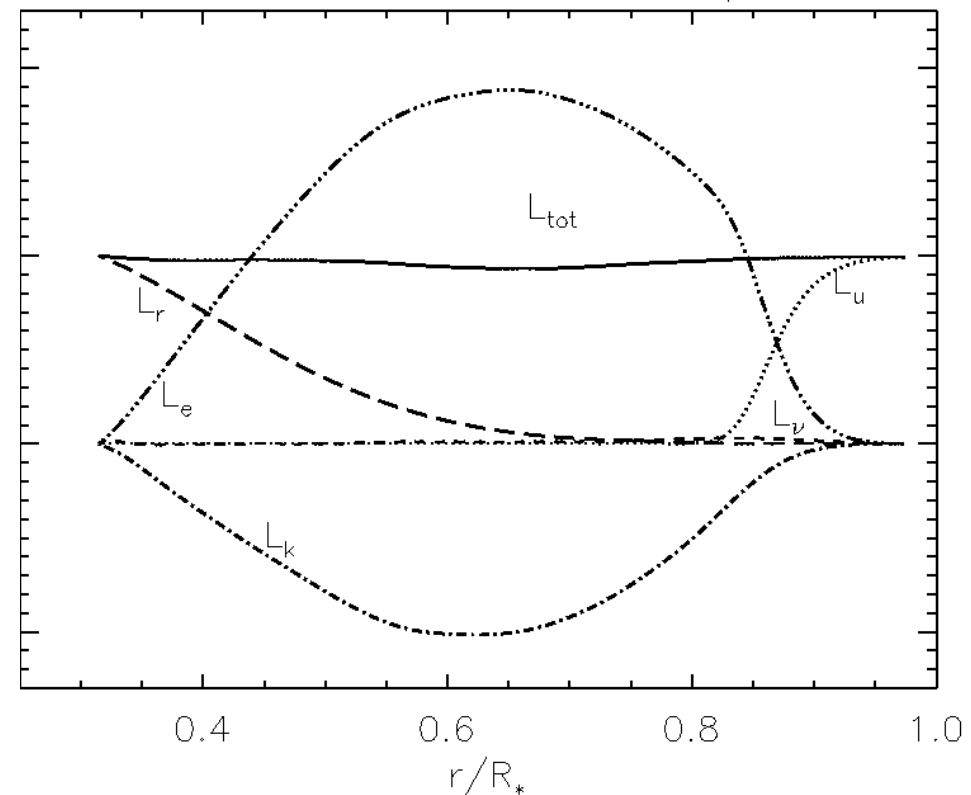
→ differs from MLT predictions

Negative kinetic energy represents up to 100% of the flux in the inner part of the domain

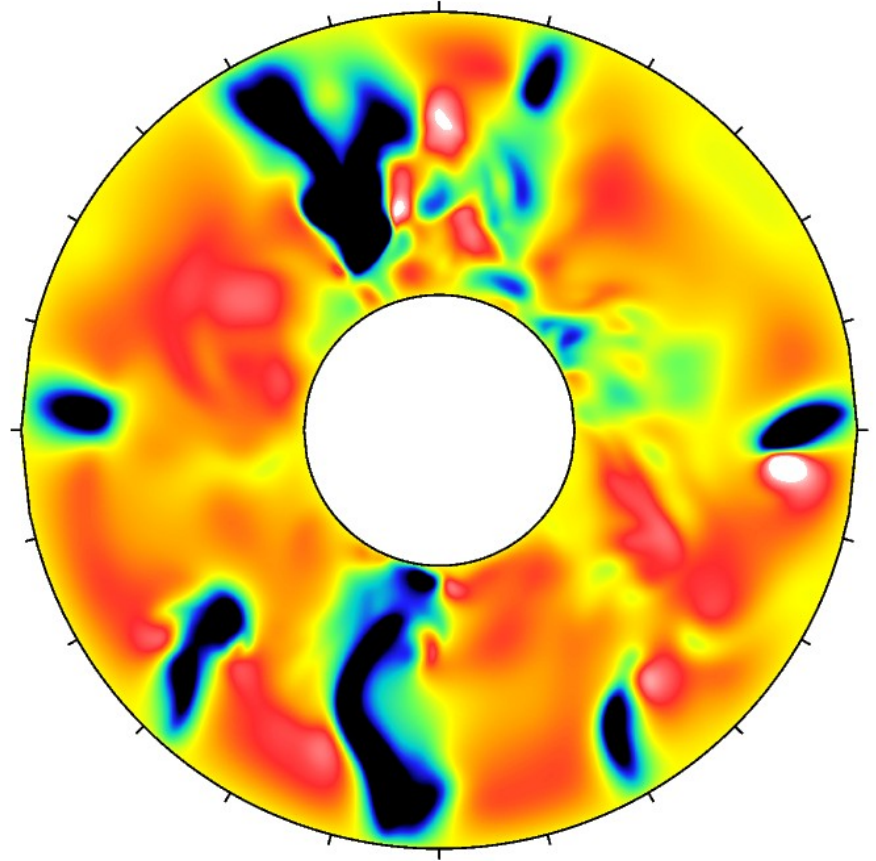
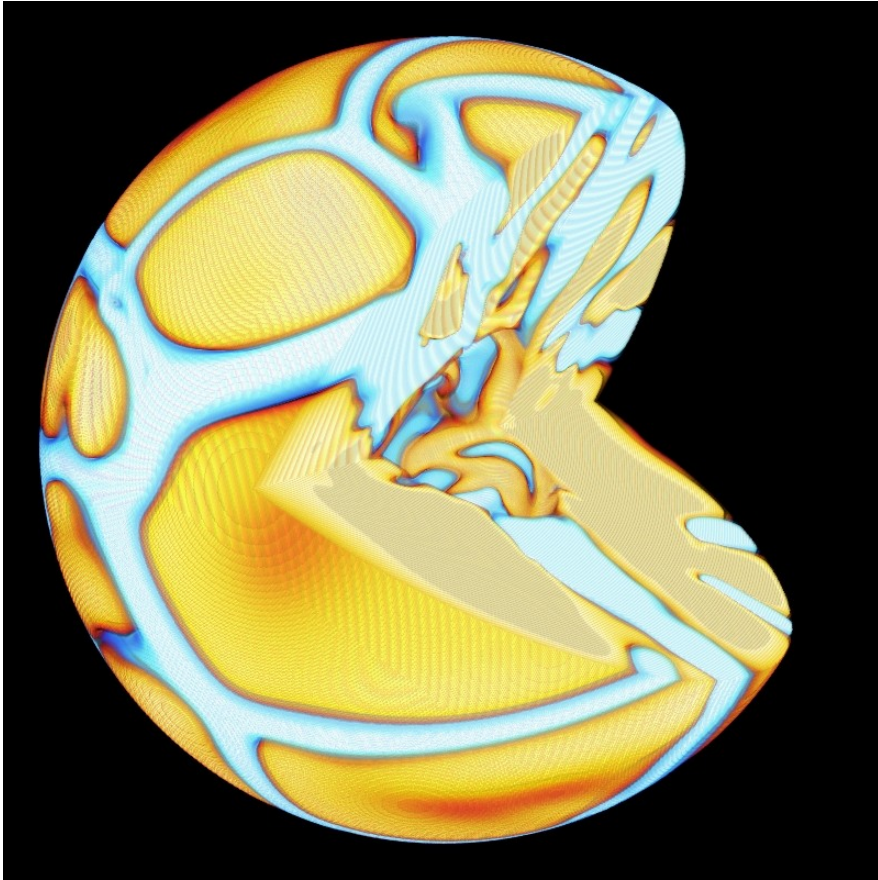
POLLUX, slow rotation, multipole



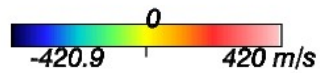
POLLUX, slow rotation, dipole



Convective motions

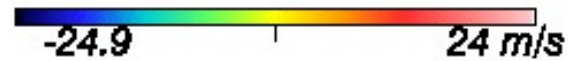
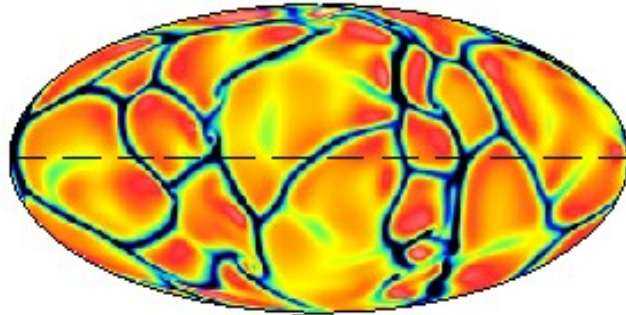


V_r

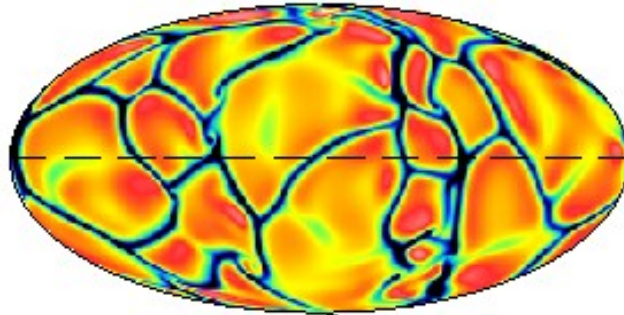


Convective motions – Radial velocity fluctuations

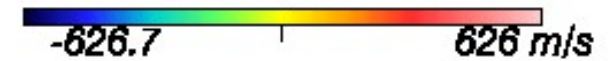
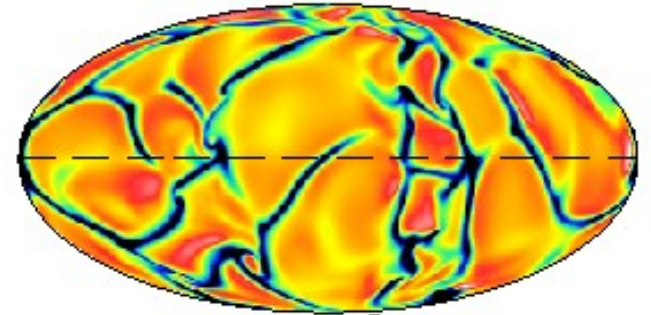
$V_r @ r = 8.8 R_\odot$



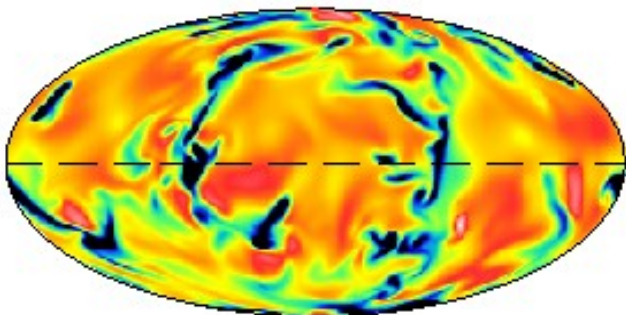
$V_r @ r = 8.5 R_\odot$



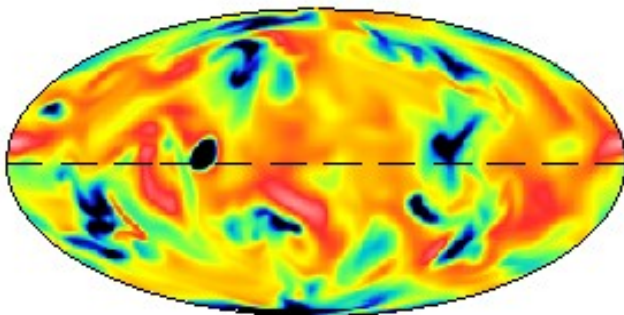
$V_r @ r = 7.4 R_\odot$



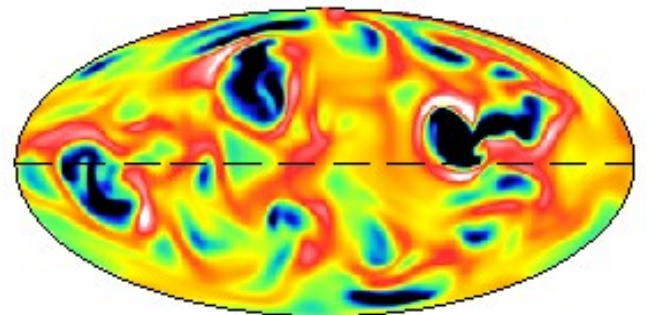
$V_r @ r = 5.8 R_\odot$



$V_r @ r = 4.1 R_\odot$



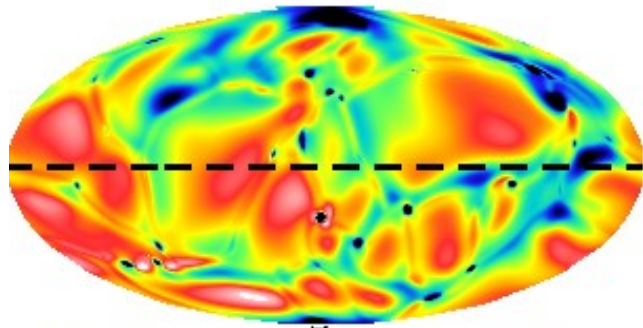
$V_r @ r = 3.1 R_\odot$



Case A – Similar pattern and values for case B

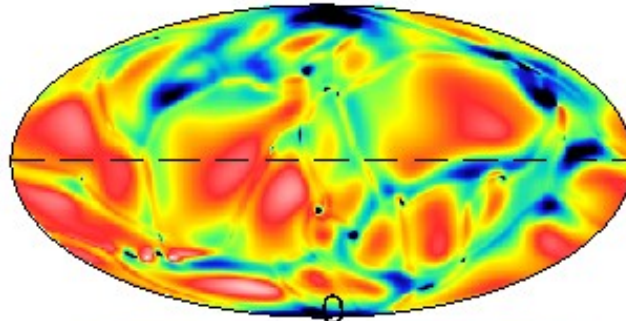
Convective motions – Temperature fluctuations

$T @ r = 8.8 R_{\odot}$



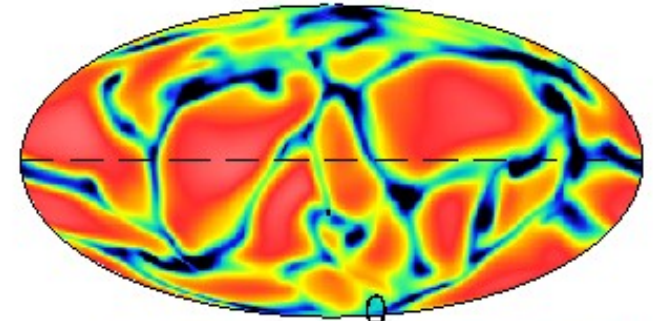
-10.9 10 deg K

$T @ r = 8.5 R_{\odot}$



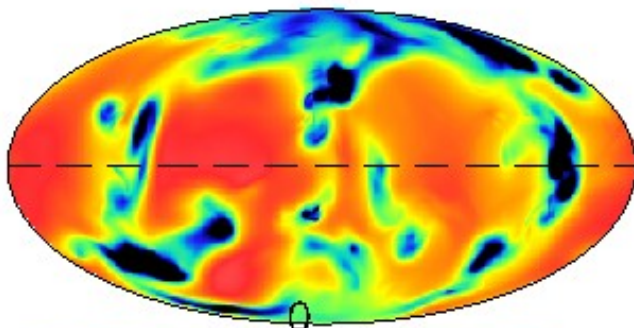
-10.9 10 deg K

$T @ r = 7.4 R_{\odot}$



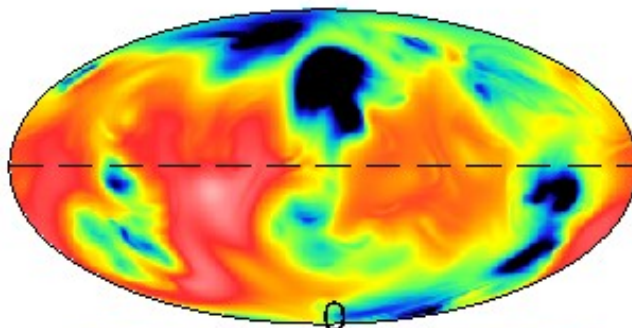
-16.3 16 de

$T @ r = 5.8 R_{\odot}$



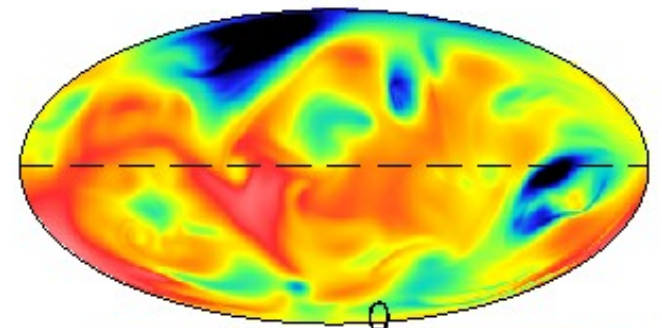
-11.8 11 deg K

$T @ r = 4.1 R_{\odot}$



-6.4 6 deg K

$T @ r = 3.1 R_{\odot}$

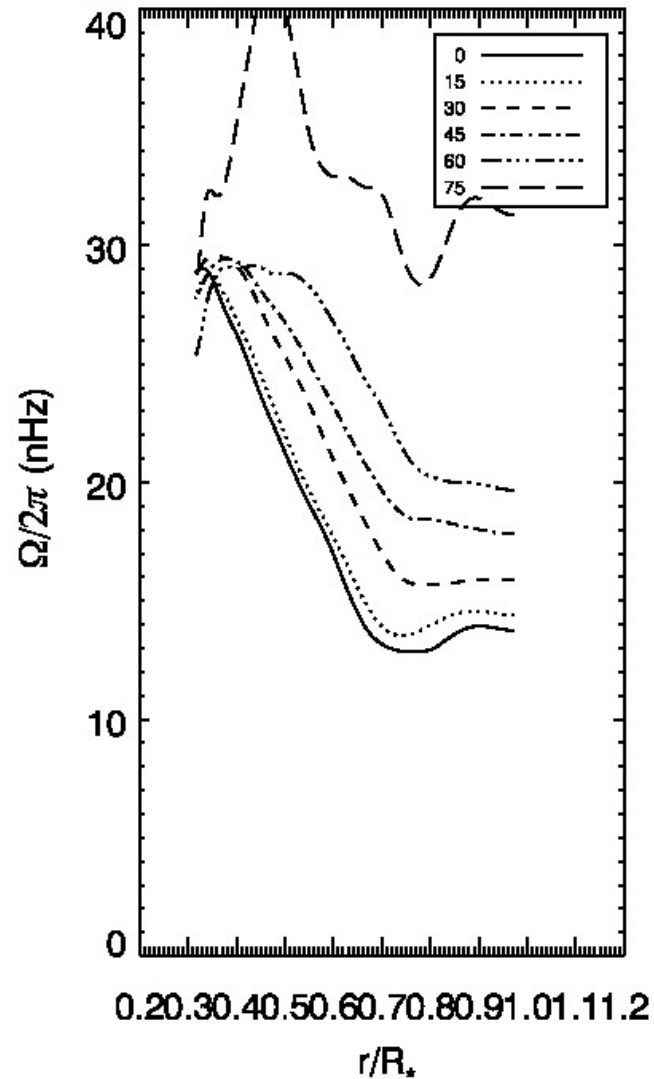
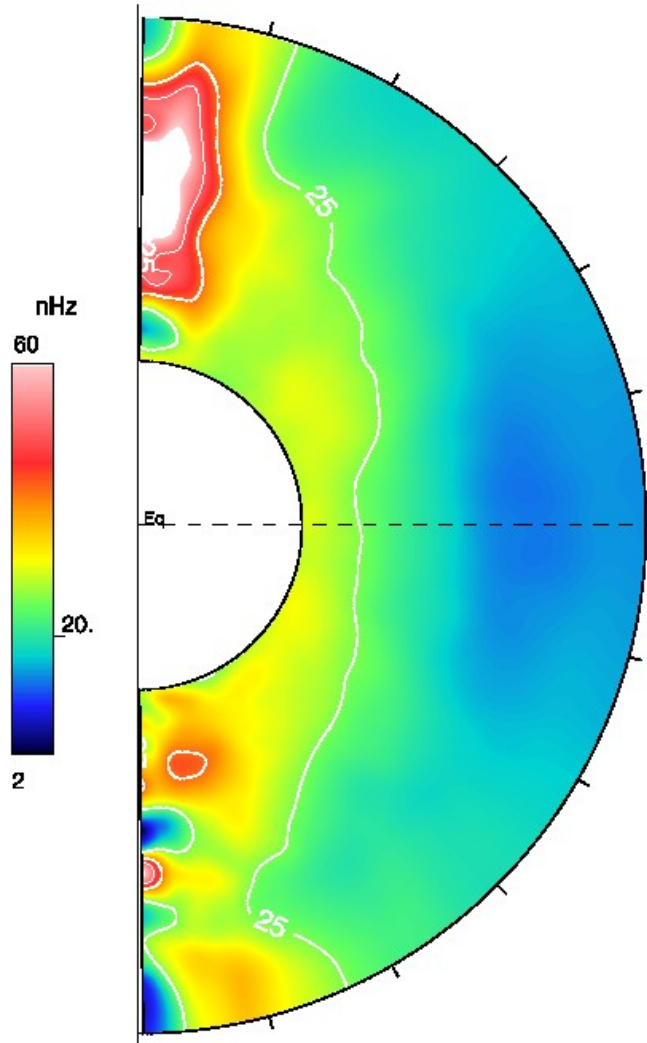


-5.6 5 deg

Case B – Similar pattern and values for case A

Rotation

Moderate differential rotation in the radial direction, moderate in latitude.

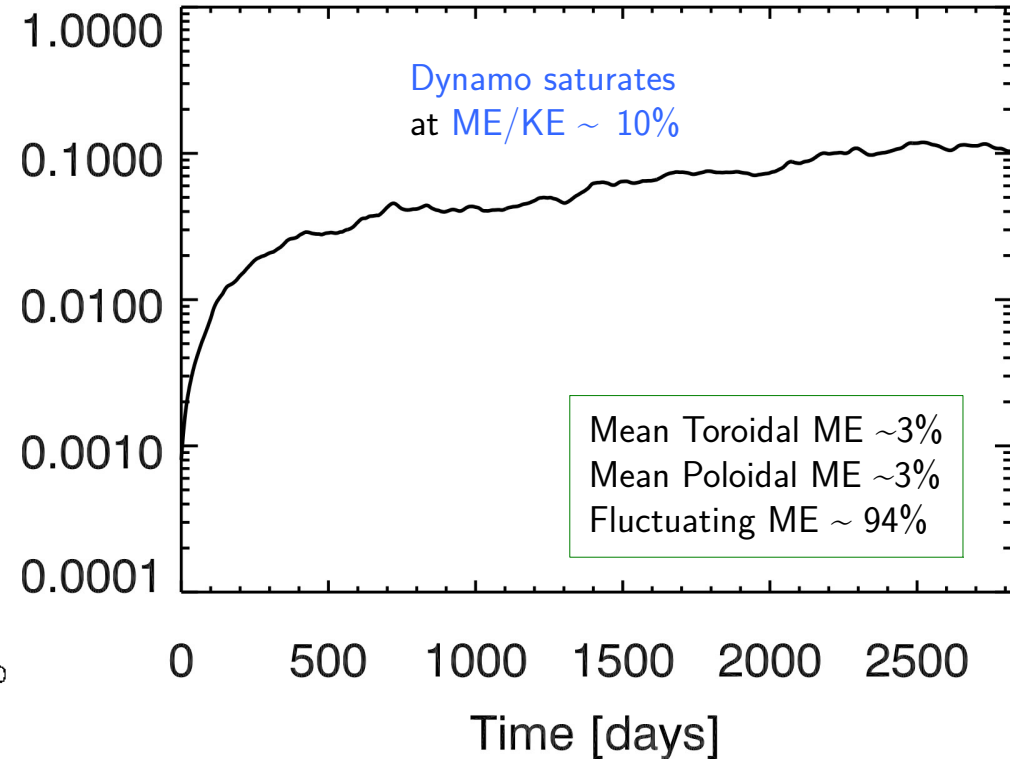
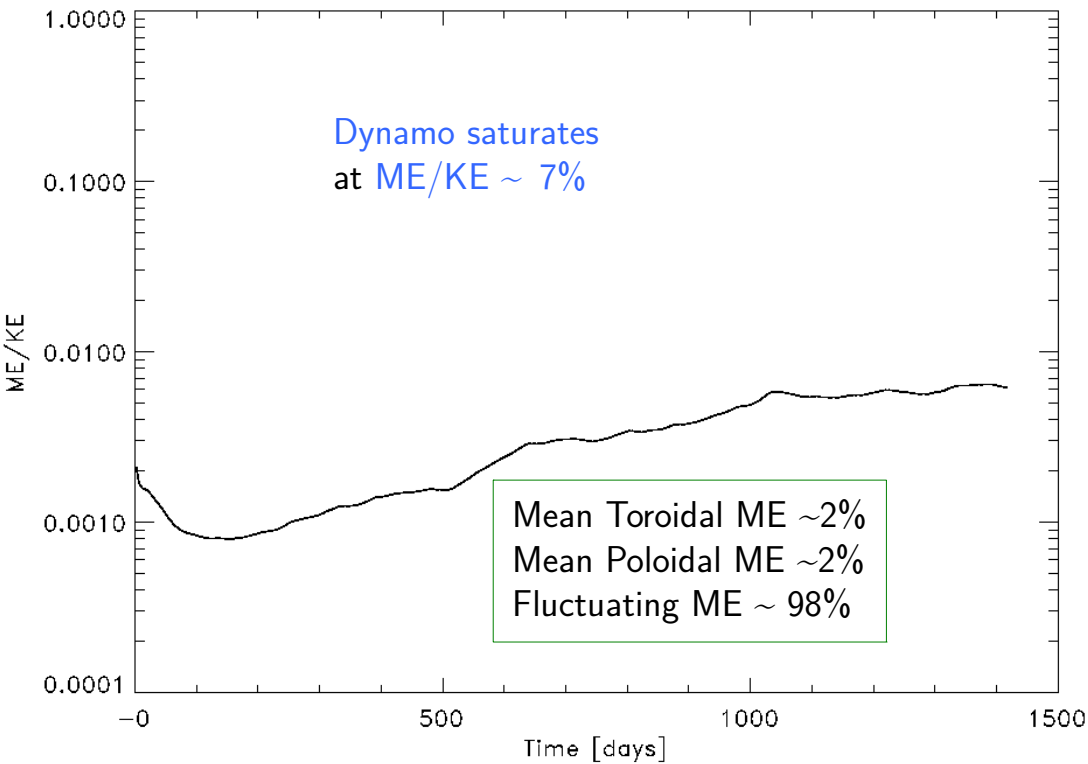


Mean over 1 P_{rot} – Case A

Dynamo action starting with seeds of different topologies

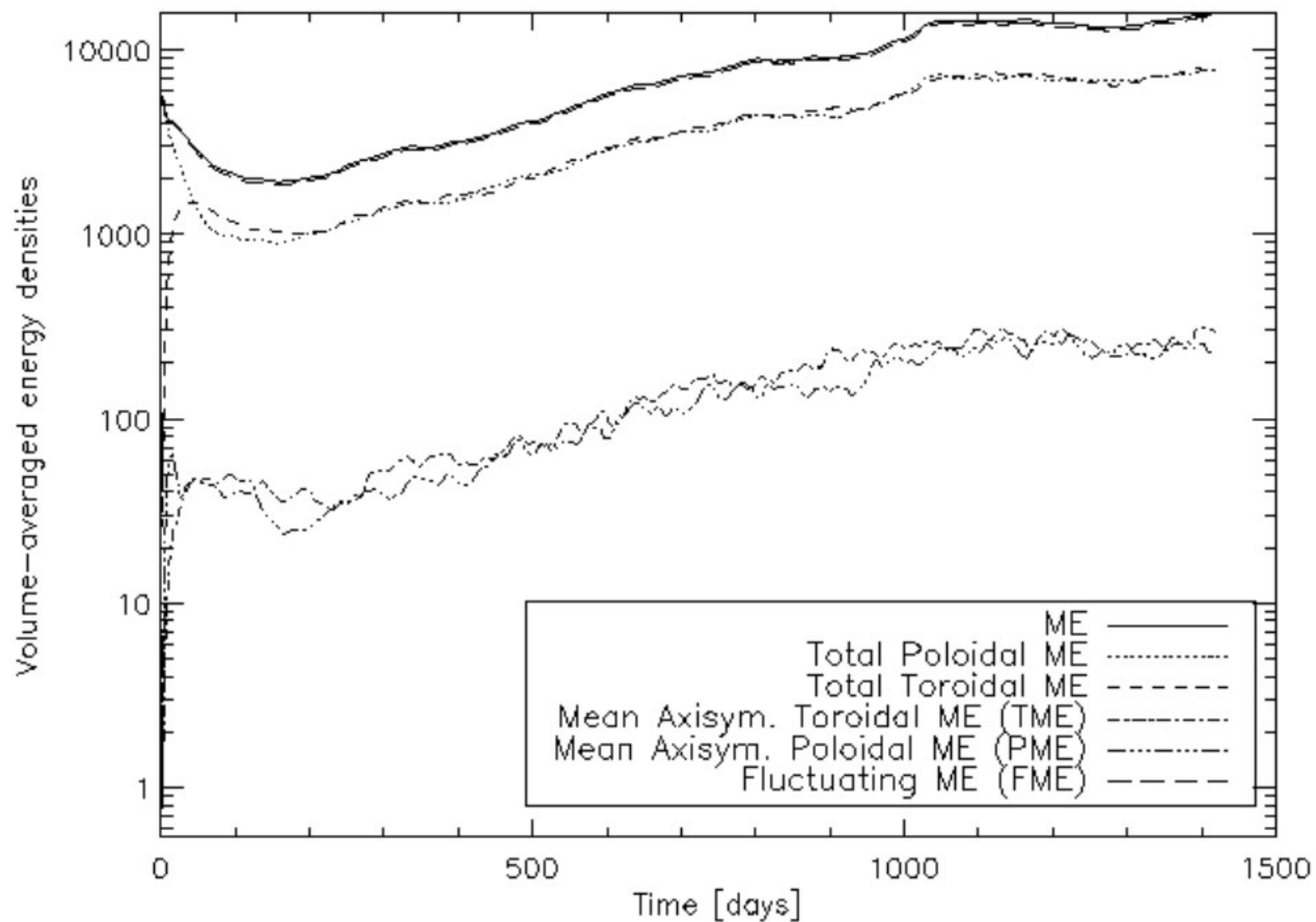
Case A : seed multipolar field

Case B : seed dipolar field

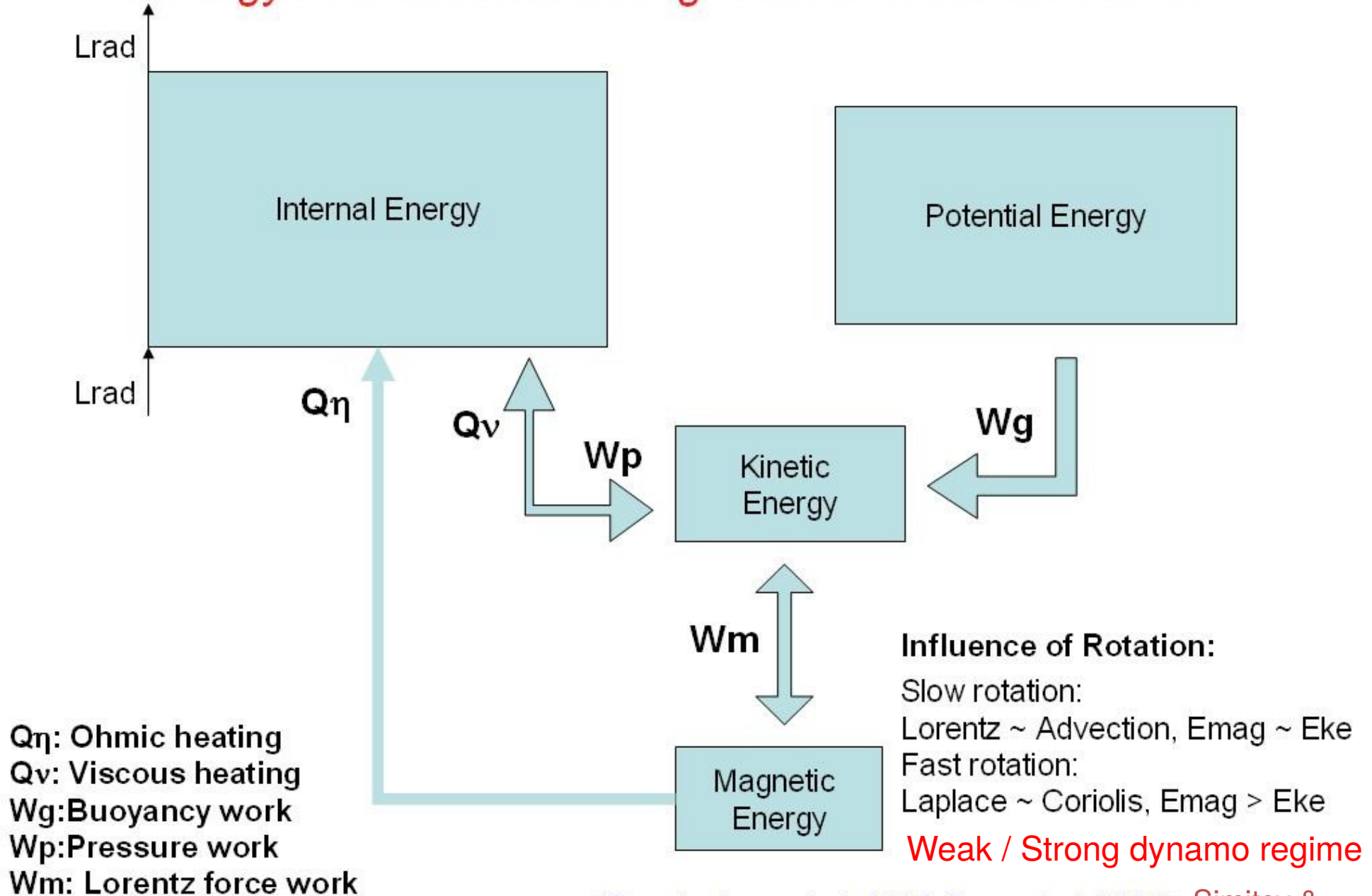


Initial $ME/KE < 0.1\%$

Evolution of magnetic energies in simulation A

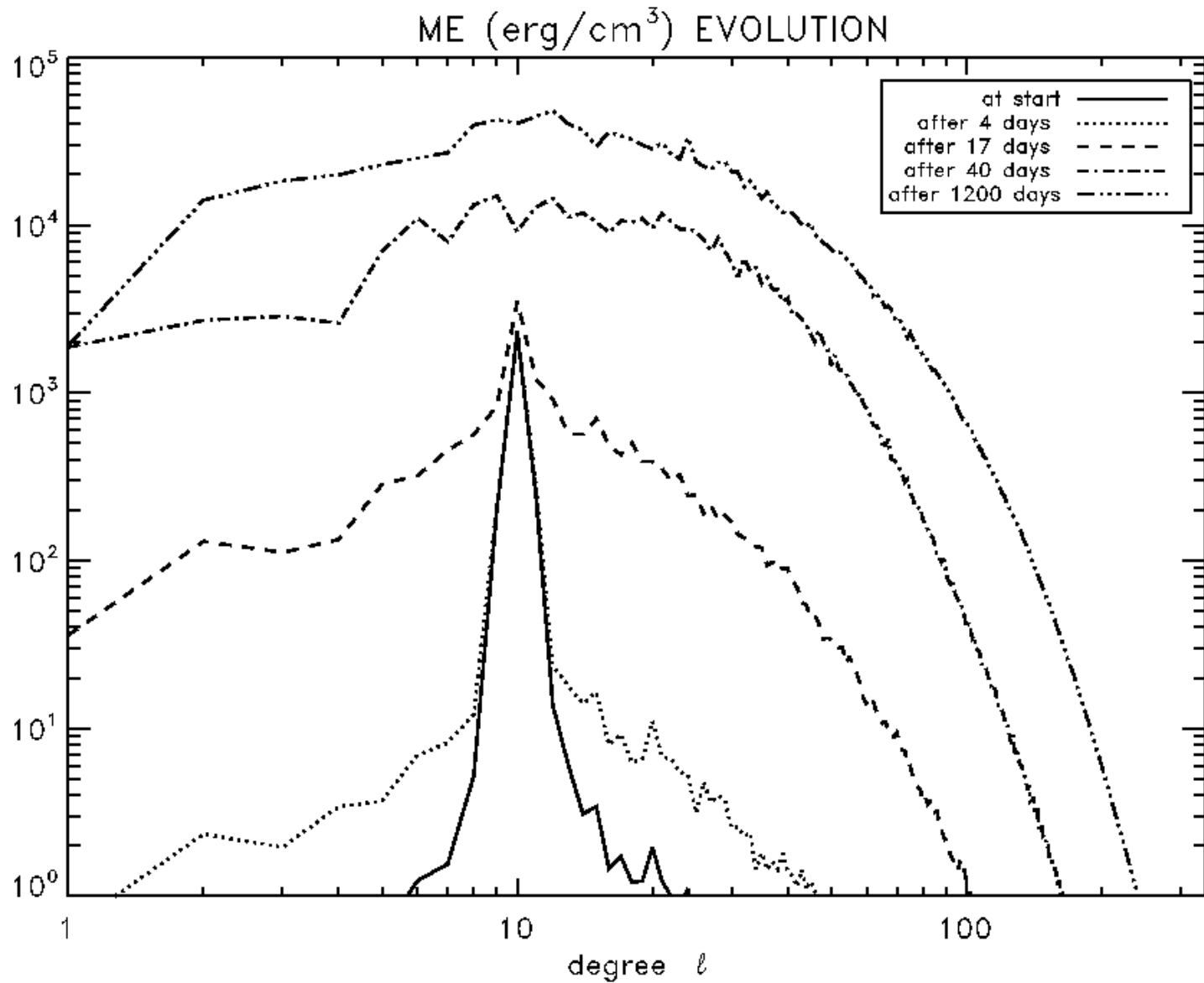


Energy Reservoirs in a Magnetized Convection Zone



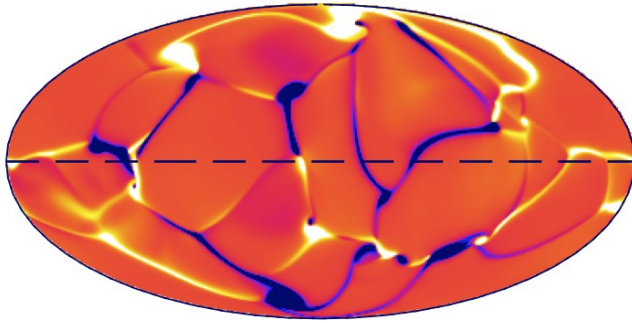
(Brandenburg et al. 1996, Brun et al. 2004) Simatev & Busse 2009

Evolution of magnetic energy spectrum in case A

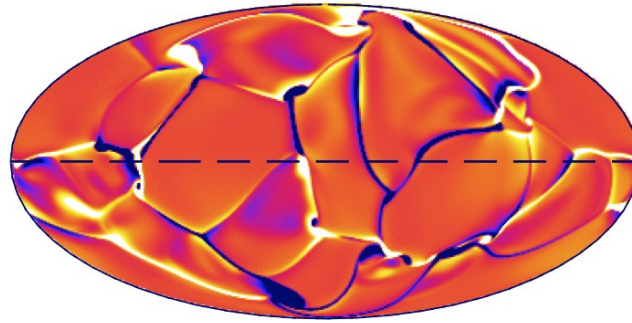


Magnetic field – radial component B_r

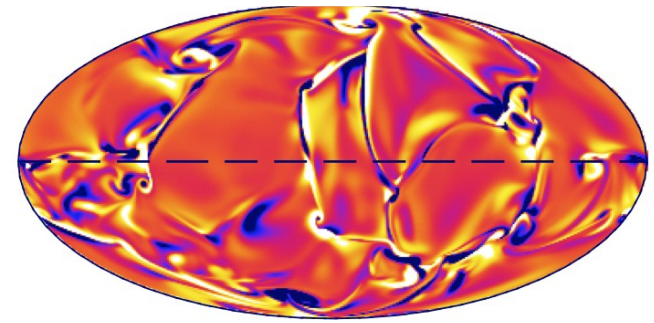
$B_r @ r = 8.8 R_\odot$



$B_r @ r = 8.5 R_\odot$

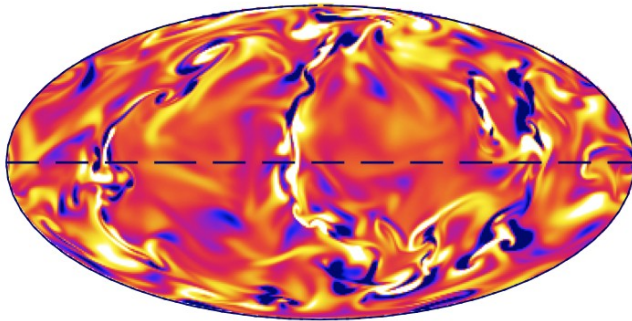


$B_r @ r = 7.4 R_\odot$

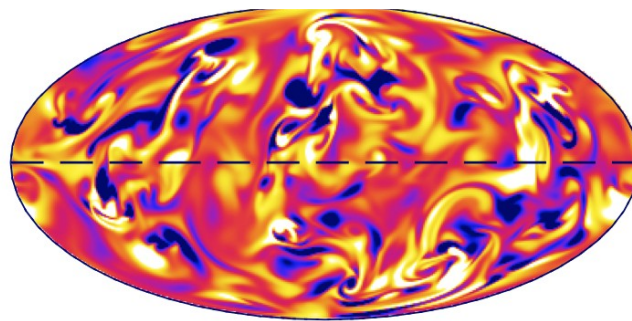


case B

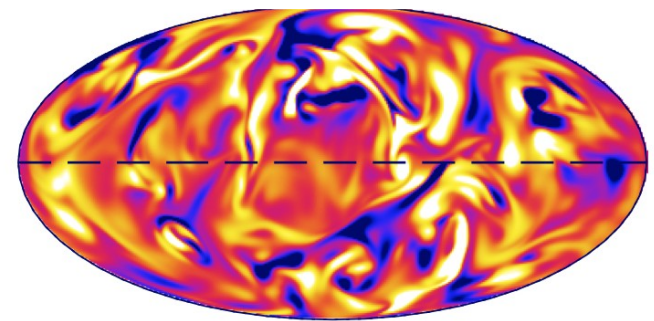
$B_r @ r = 5.8 R_\odot$



$B_r @ r = 4.1 R_\odot$

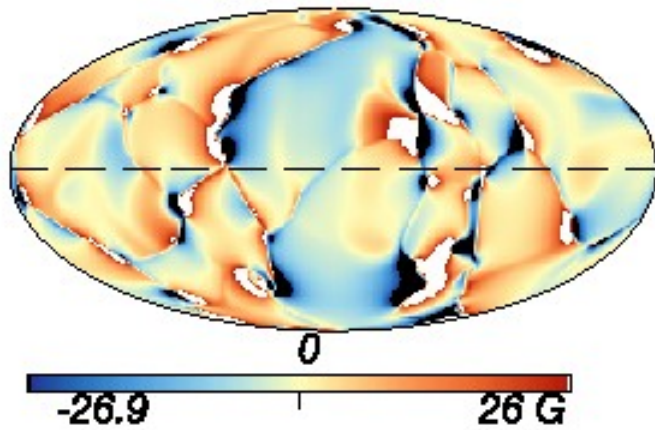


$B_r @ r = 3.1 R_\odot$

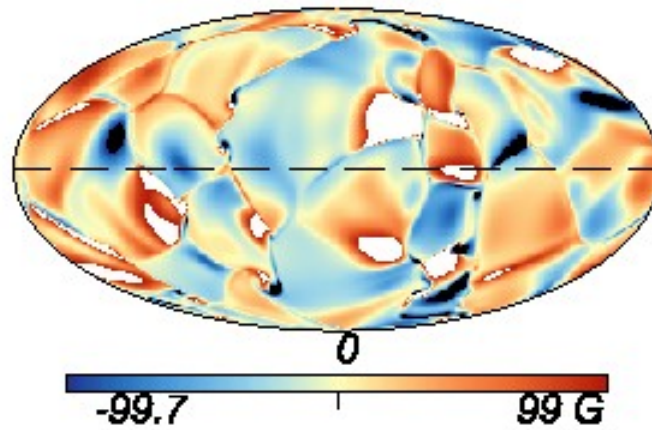


Magnetic field – azimuthal component B_ϕ

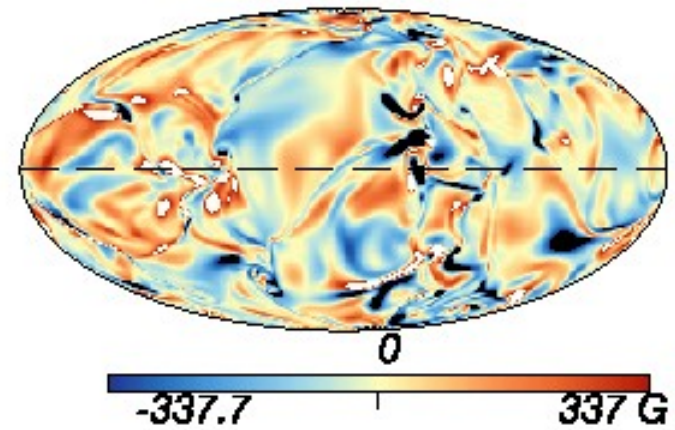
$B_\phi @ r = 8.8 R_\odot$



$B_\phi @ r = 8.5 R_\odot$

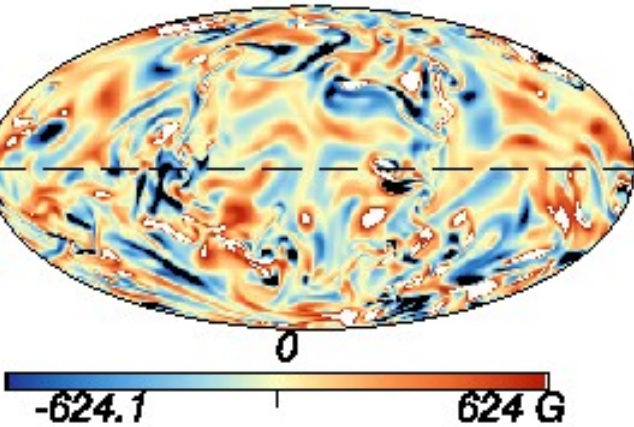


$B_\phi @ r = 7.4 R_\odot$

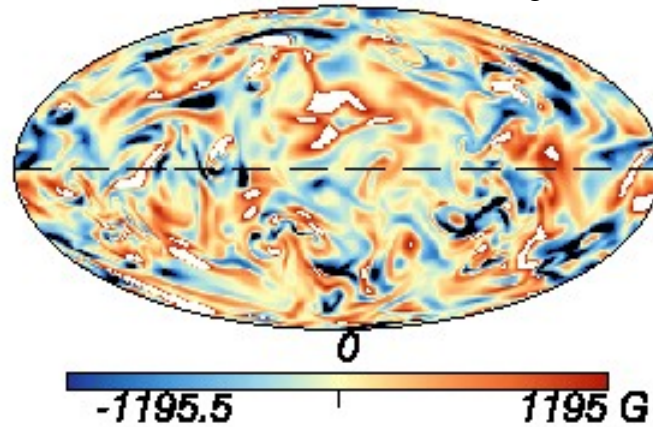


case A

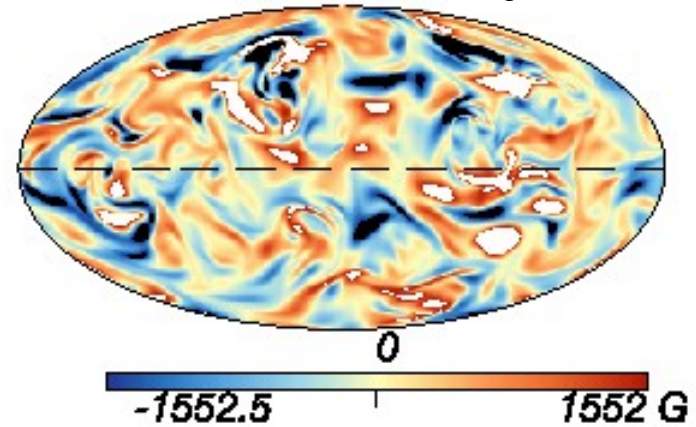
$B_\phi @ r = 5.8 R_\odot$



$B_\phi @ r = 4.1 R_\odot$

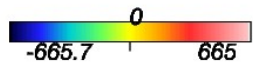
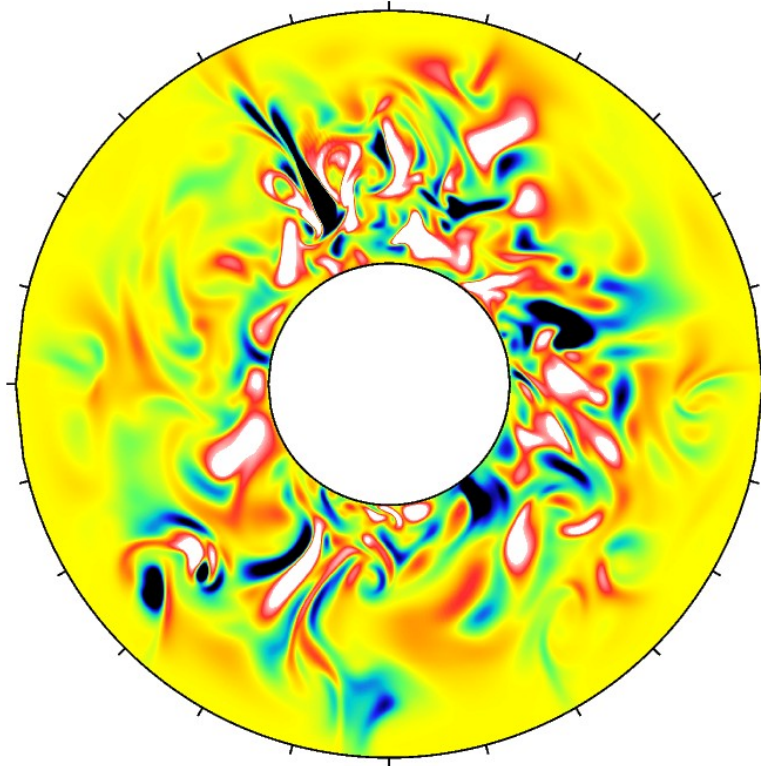


$B_\phi @ r = 3.1 R_\odot$

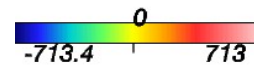
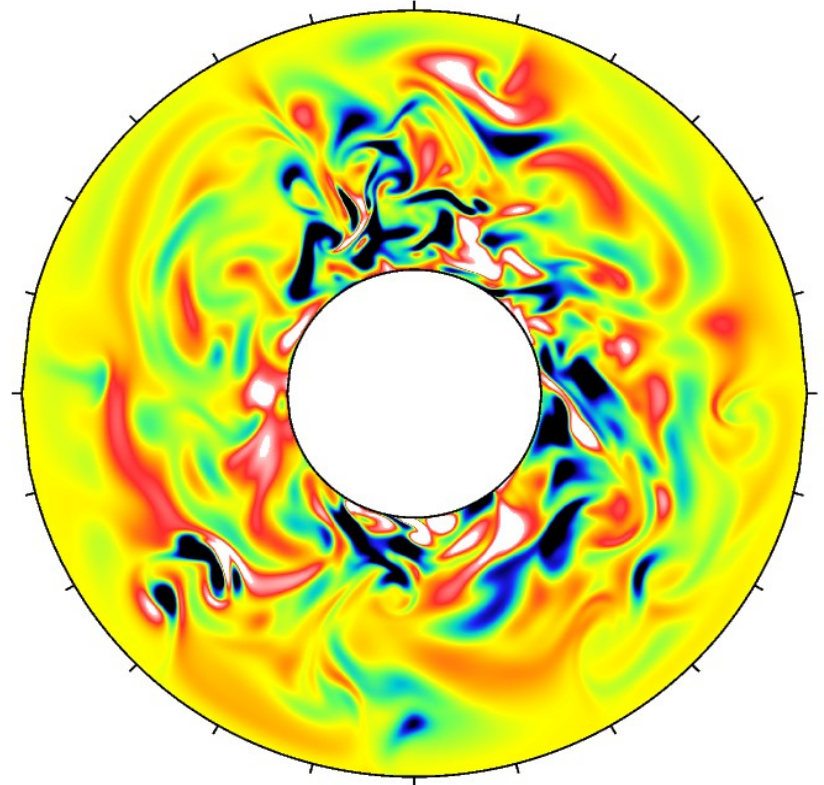


Magnetic field – radial and azimuthal euqatorial slices

Case A



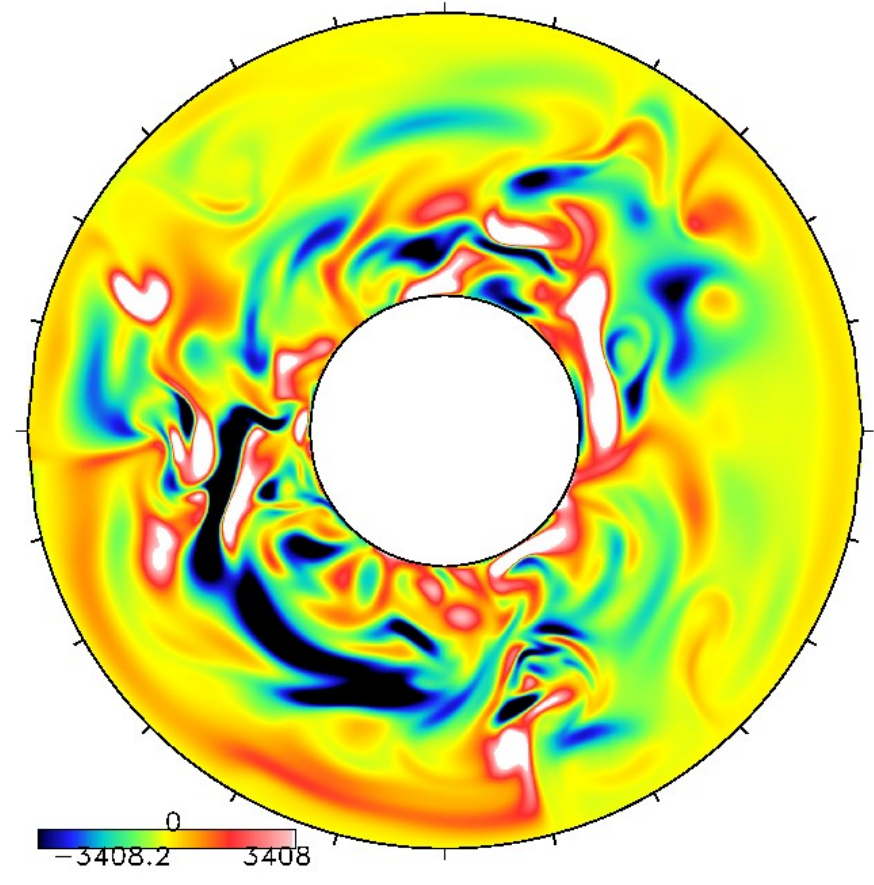
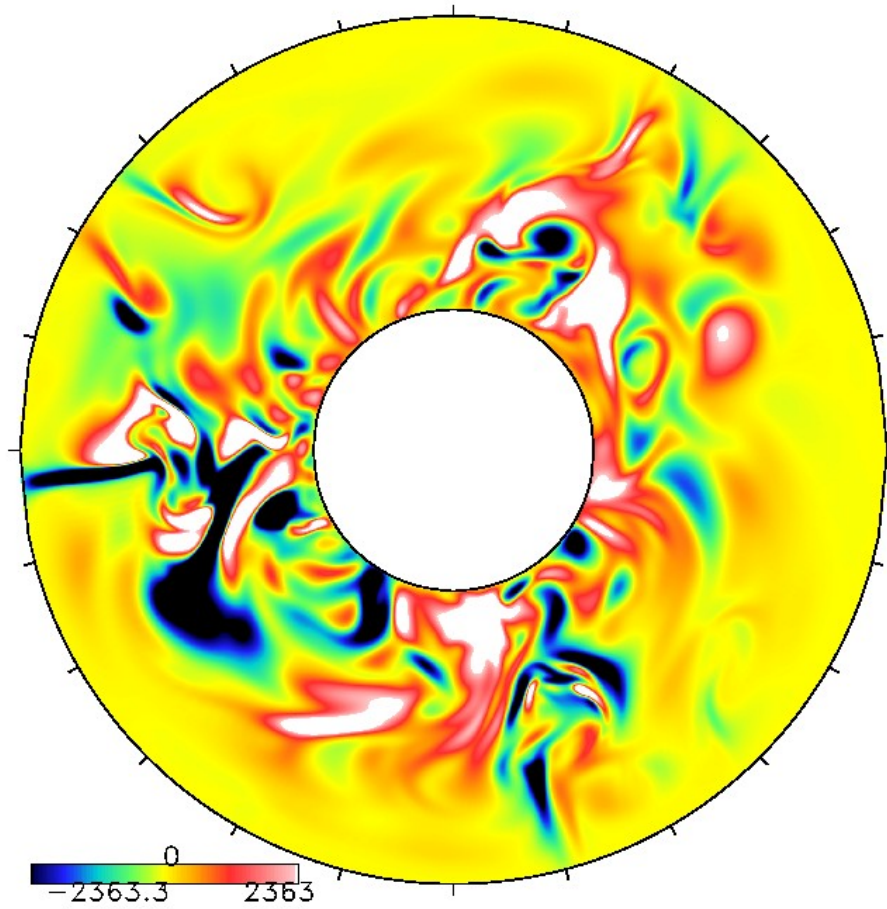
B_r



B_ϕ

Magnetic field – radial and azimuthal equatorial slices

Case B

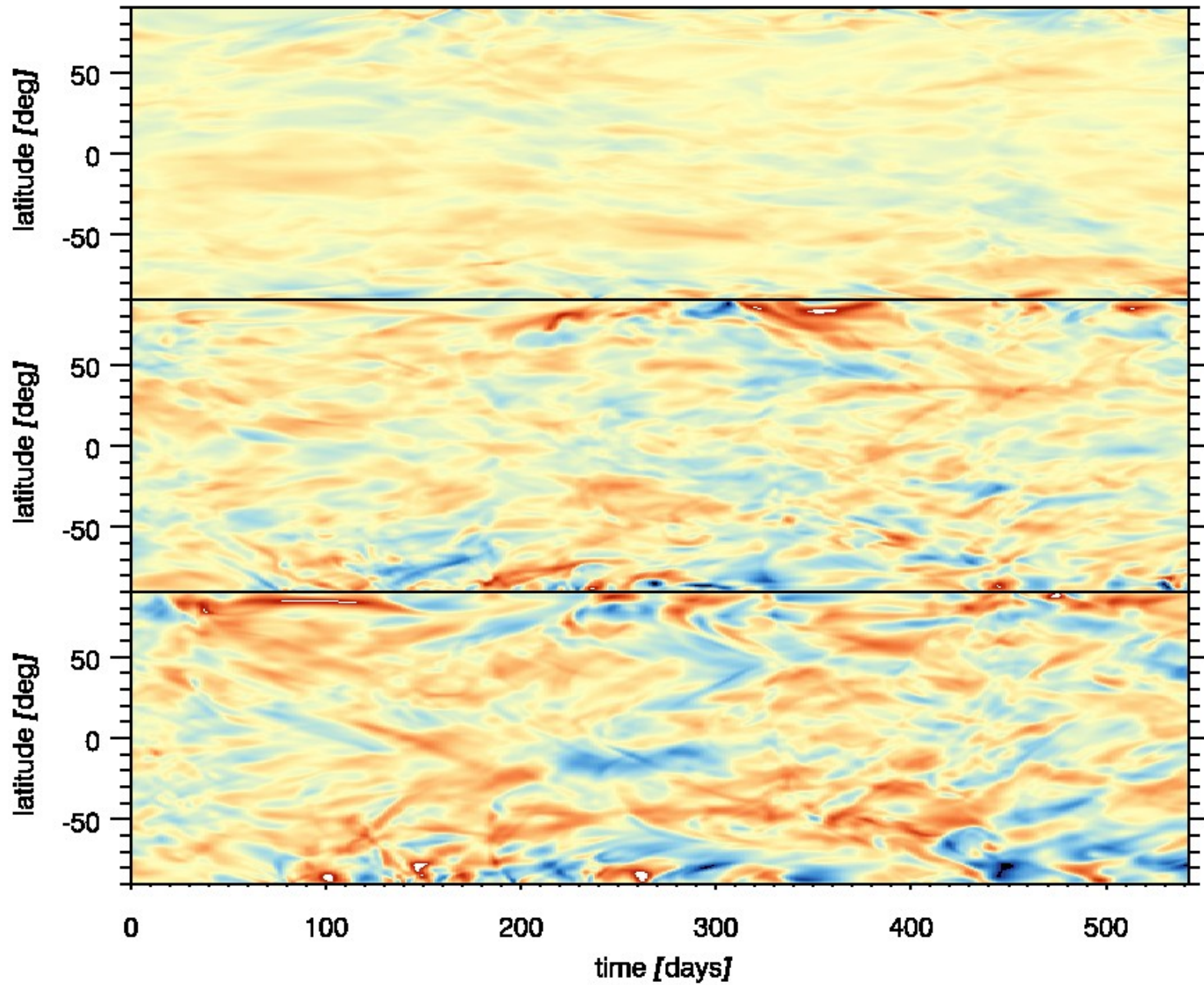


B_r

B_ϕ

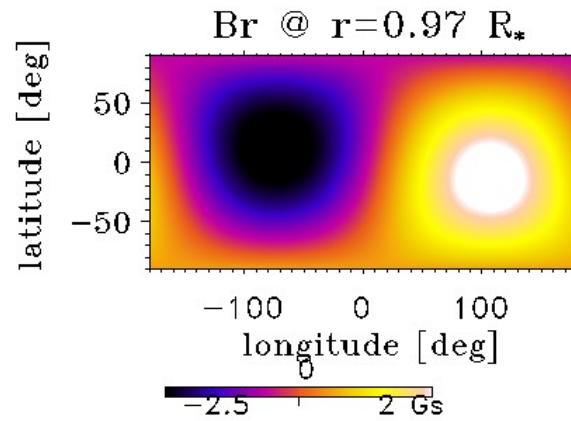
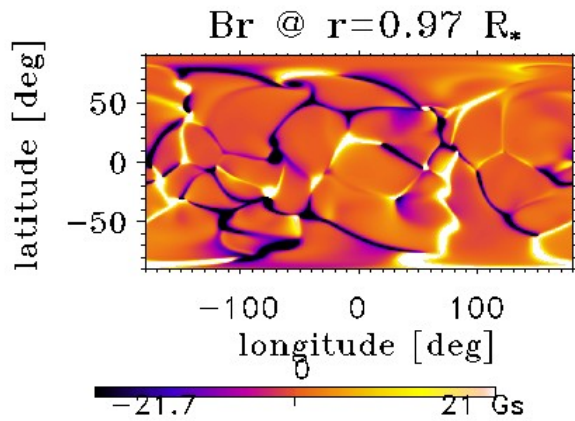
Evolution of azimuthal magnetic field in case A

Colour scale ± 200 G

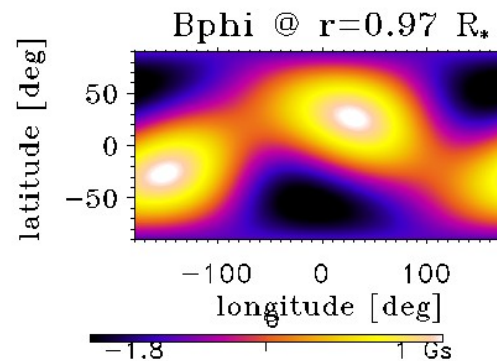
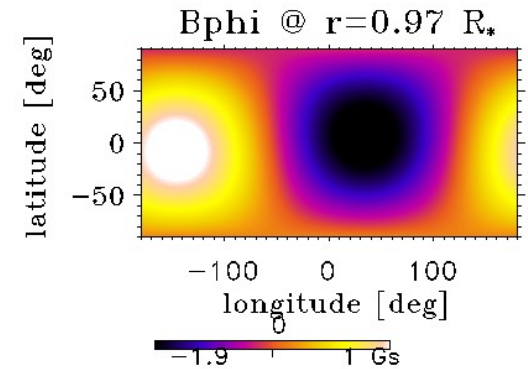
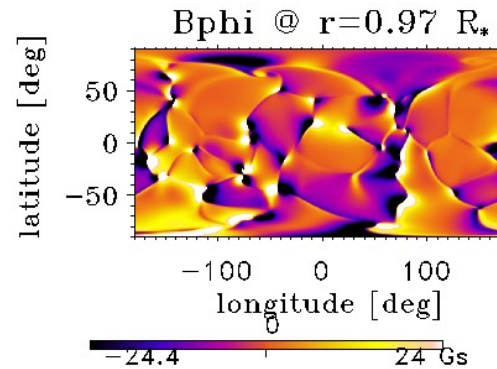
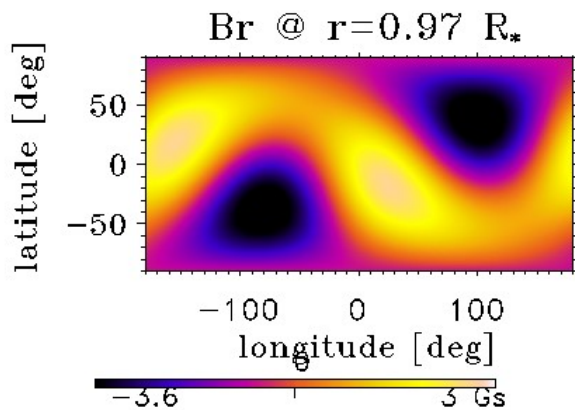


Magnetic field intensity in the dipole and quadrupole

Case A



Radial field

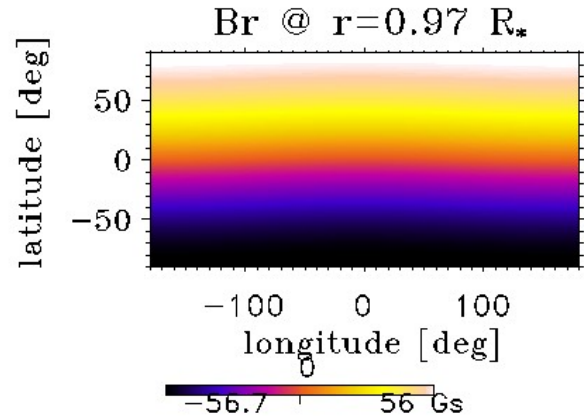
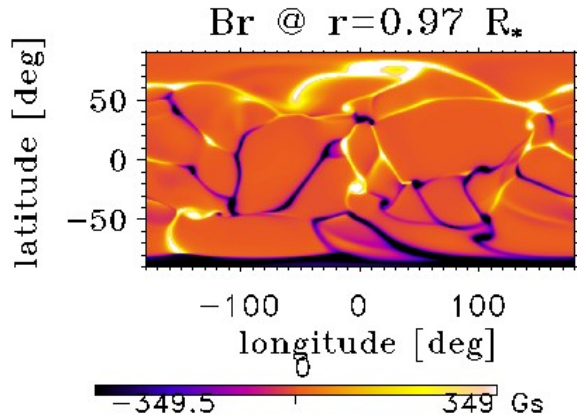


Longitudinal field

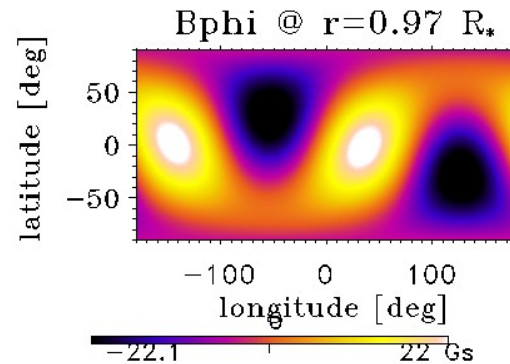
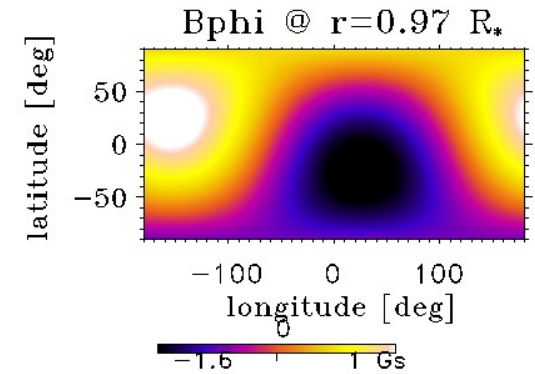
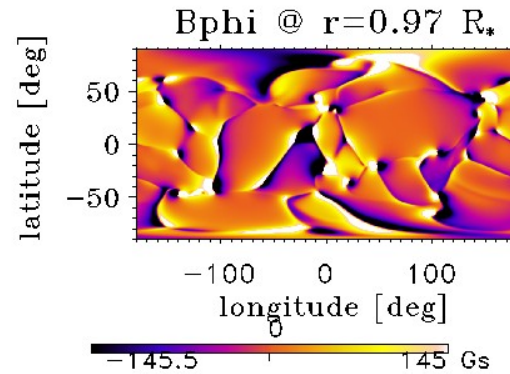
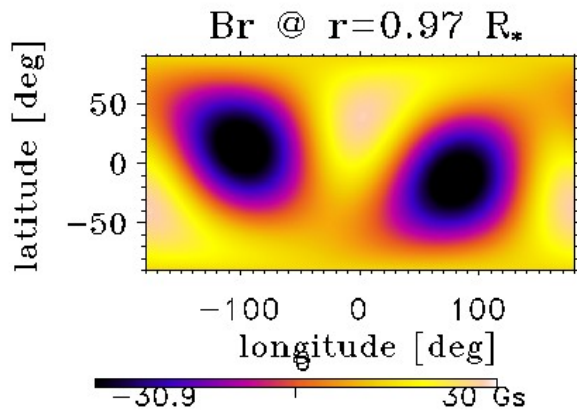
The intensity of the toroidal field in the dipole and quadrupole is compatible with the value of $|B_t|_{\max}$ derived from Stokes V of about 1 Gauss

Magnetic field intensity in the dipole and quadrupole

Case B



Radial field



Longitudinal field

The intensity of the toroidal field in the quadrupole is much larger for case B (about 20 G) and thus one order of magnitude higher than the $|B_l|_{\max}$ derived from spectropolarimetry

Take away messages

Slowly rotating RGB star can easily **generate magnetic field via dynamo action** in their convective envelope

They have some **differential rotation mostly radial**

Overall field energy is around **10% of kinetic energy**

The dipolar and multipolar seeds lead to a dynamo possessing a **dipolar mode compatible with that observed**

Parameter study will tell us how sensitive it is to initial conditions (A vs A_p star)



Various Dynamo Regimes and Scalings

Equilibrium field : $B_{\text{eq}} \sim \text{sqrt}(8\pi P_{\text{gaz}}) \sim \text{sqrt}(\rho_*)$

If magnetic Reynolds number $Rm \sim 1$, $v = \eta/L$, then

Laminar (weak) scaling: Lorentz \sim diffusion \rightarrow

$$B_{\text{weak}}^2 \sim \rho v \eta / L^2$$

Turbulent (equipartition) scaling: Lorentz \sim advection \rightarrow

$$B_{\text{turb}}^2 \sim \rho v^2 \sim \rho \eta^2 / L^2 \Leftrightarrow |B_{\text{weak}}| \sim |B_{\text{turb}}| P_m^{1/2}$$

Magnetostrophic (strong) scaling: Lorentz \sim Coriolis \rightarrow

$$B_{\text{strong}}^2 \sim \rho \Omega \eta$$

With ρ density, ν kinematic viscosity, η magnetic diffusivity, Ω rotation rate, v , L characteristic velocity & length scales, $P_m = \nu/\eta$ the magnetic Prandtl number